

Residual circulation and thermohaline distribution of the Ría de Vigo: a 3-D hydrodynamical model*

S. TORRES LÓPEZ¹, R.A. VARELA¹ and E. DELHEZ²

¹GOFUVI, Facultad de Ciencias, Universidade de Vigo, Campus de Lagoas-Marcosende, 36200 Vigo, Spain.

²G.H.E.R., University of Liège, B5 Sart Tilman, 4000 Liège, Belgium.

SUMMARY: A three-dimensional, non-linear, baroclinic model is described and tested for the first time to study the residual circulation and the thermohaline distribution of the Ría de Vigo (NW Spain) at short time scales and under different wind stress regimes. Two markedly different realistic scenarios were chosen: northerly upwelling-favourable winds and southerly downwelling-favourable winds. The numerical experiments carried out indicate that the hydrodynamic regime of the Ría de Vigo is mostly a consequence of wind events. As could be expected, moderately strong North winds reinforce the normal (positive) Ría circulation, while winds blowing from the South, when sufficiently strong, reverse the typical circulation pattern and reduce the characteristic outgoing velocities and the flushing time inside the Ría. The temperature and salinity fields generated by the model in both situations were compared with observations and found to be in qualitatively good agreement, supporting the 3D velocity field distribution.

Key words: numerical models, ría de Vigo, upwelling-downwelling, thermohaline distributions, currents.

INTRODUCTION

The coast of Galicia (NW Spain) is characterised by the presence of several embayments, usually called "Rías". Significant human settlements are associated with these embayments, which are also among the most important mussel culture areas of the world. In particular, the Ría de Vigo, located near the boundary between Portugal and Spain between (42° 09'N, 8° 36'W) and (42°21'N, 8°54'W) supports the highest population among the Rías (more than 500,000) and is probably the best known of all of them from an oceanographic point of view (Fig. 1).

The Ría de Vigo freshwater input comes from the Verdugo-Oitabén river (about 79% of the total annu-

al mean freshwater flow to the Ría, López-Jurado, 1985), from the Lagares river (~19%, Alejo, 1994) and also partially from the city of Vigo (about 2%, *c.f.* Fraga, 1996). The volume of freshwater flowing into the Ría is small (the total volume of freshwater present in the Ría represents only 2% (Fraga, 1996) of the total volume of the Ría) and much smaller than the water exchanged with the shelf and the ocean. It will be shown, however, to have a non-negligible influence on the residual flow in the Ría.

The local winds, generally blowing along the axis of the Ría (because of the particular orography), can enhance or reduce the water flushing time (between 5 and 7 days) inside the Ría. The mixing between surface waters (highly dependent on local meteorological conditions) and the presence of North Atlantic Central Water (NACW), determines

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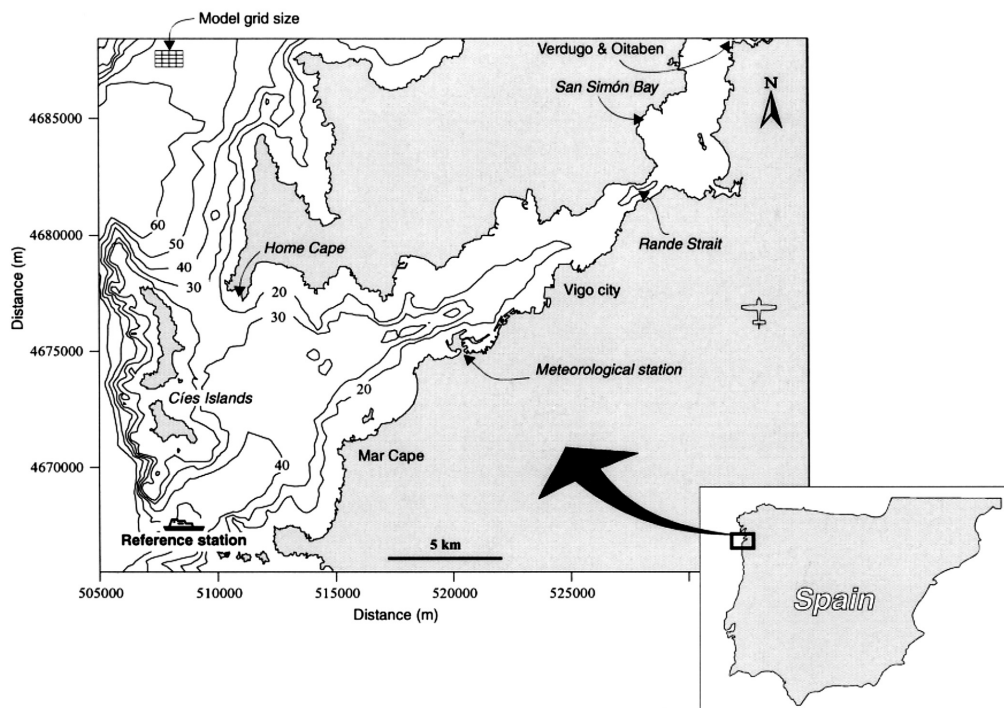


FIG. 1. – Map of the Ría de Vigo (NW Spain), with the location of the meteorological station, the reference station where boundary conditions profiles were measured and a scale grid. Contour lines represent depth in meters.

the characteristics of the oceanic waters flowing into the Ría.

The synoptic-scale meteorological field on the Galician shelf is dominated by upwelling-favourable winds (north, northwest direction) from early spring to the end of the summer. This seasonal wind field is, however, intermittently altered by relaxation or downwelling-favourable winds (south, southwest direction). Transitions between upwelling and downwelling-favourable winds tend to occur by September and October (Blanton *et al.*, 1984; Bakun and Nelson, 1991, Castro *et al.*, 1994). Although several studies describe the strong influence of the wind variability on the water mass dynamics in the Rías Baixas (Tilstone *et al.*, 1994; Rosón *et al.*, 1995; Doval *et al.*, 1997; Rosón *et al.*, 1997), no hydrodynamic 3-D model has yet been applied to study these influences.

Margalef and Andreu (1958) first estimated vertical velocities in the Ría de Vigo, giving mean annual values in the order of a few (1-2) metres per day. Later, using different box models, Prego (1988), Prego and Fraga (1992) and Ríos *et al.* (1992) calculated the horizontal residual velocities and the vertical associated mixing for a one-year period. Prego (1988), who divided the Ría into 5 horizontal sections of 2 vertical boxes, each separated by the halocline/pycnocline, obtained horizontal

residual velocities (for a given date) ranging from 0.05 m s^{-1} near the Verdugo-Oitabén River to about 0.01 m s^{-1} next to the oceanic border.

Nogueira *et al.* (1997a; 1997b; 1998) used a statistical approach to model the hydrography of the Ría de Vigo using historical data collected from 1987 on a fortnightly basis, at a fixed station located in its central part. The temporal evolution of the thermohaline properties, nutrients and chlorophyll *a* were fitted to a linear stochastic model. The results were used to describe the dynamics of the system, its seasonal patterns and its long-term trends. It was shown that, when helped by appropriate meteorological data, Box-Jenkins Transfer Function models adequately describe the dynamics of the Ría de Vigo ecosystem and may be useful as a first approximation to developing more sophisticated *deterministic* models.

Further developments in the modelling of the Ría de Vigo were made by Taboada *et al.* (1998), who developed a 3-D baroclinic model to study the seasonal variations of the Ría de Vigo's residual circulation. Their model used a very rough turbulence closure scheme (i.e. constant and uniform eddy diffusion and eddy viscosity coefficients), taking a constant horizontal grid spacing and eight vertical layers with a double- σ vertical coordinate system. Since the authors developed the model with long-term analysis

in mind, the wind stress was argued to be zero. This circulation model was compared to and found to be in reasonable agreement (order of magnitude) with the box model of Prego and Fraga (1992).

Very recently, Gómez-Gesteira *et al.* (1999) used the model of Taboada *et al.* (1998) in a 2D (i.e. vertically averaged) mode to track the dispersion of passive pollutants in the Rías of Vigo and La Coruña. They compared the residence time of particles within both harbours, concluding that the Ría de Vigo has a smaller residence time, probably due to the interference of a dock in La Coruña harbour.

The objective of this paper is to describe a full baroclinic 3D-circulation model of the Ría de Vigo that overcomes the limitations of earlier model attempts. The proposed model has a refined turbulence closure scheme and could be easily coupled to biogeochemical modules. It is used to evaluate the general influence of the meteorological forcing on the hydrodynamic conditions in the Ría de Vigo. Even when the model uses the main harmonic tidal components to simulate the tide inside the Ría de Vigo, no special tidal current analysis is included in this paper in order to focus on the residual circulation, which is mainly settled by wind forcing.

To achieve this goal, two groups of model simulations were carried out, one in typical upwelling conditions and one with downwelling-favourable winds. The model results were compared with similar situations already observed and described in many papers on the Rías Baixas, focusing on the mean general circulation patterns under different wind conditions instead of on any specific, date-dependent event.

MATHEMATICAL AND NUMERICAL MODEL DESCRIPTION

Mathematical aspects

The hydrodynamic model is a modified version of the general purpose GHER 3D turbulent closure model previously applied to a wide variety of regions including the North Sea (Delhez and Martin, 1992), the Mediterranean Sea (Beckers, 1991), the Bering Sea (Nihoul *et al.*, 1989) and the Black Sea (Grégoire *et al.*, 1998). This is the first time that has been applied to a coastal estuarine domain.

The model relies on the well-known equation of the “weather of the sea” with the Boussinesq approximation (Eqs. 1 and 2, Table 1, 2). Vertical

turbulent fluxes are computed using a simple but robust and efficient closure scheme (Delhez *et al.*, 1999) based on the well-accepted evolution equation for the turbulent kinetic energy, k (per unit mass) (Eq. 4, Table 1, 2) and a parametric expression of the mixing length l (m) introduced by Nihoul (1984) (Eq. 4.3, Table 1, 2), which gives heed to the stratification and the surface wave activity.

This closure scheme has the advantage that it takes into account the advection, the diffusion, the local production-destruction and the history of turbulence through the evolution equation for k . Moreover, empirical sound velocities formulas are available for l in shallow shelf seas so that the resolution of a differential equation for the dissipation rate of k (ϵ) can be avoided and replaced by an algebraic expression of the mixing length (Nihoul, 1984) (Eq. 4.3, Table 1, 2). At the bottom, this expression respects the well-established asymptotic boundary layer. At the surface, eddies must be smaller in negligible wind conditions than when the wind buffets the surface. Therefore, the parameter $\delta \in (0.5, 1)$ is introduced to adjust the mixing length profile to wind-mixed layer conditions. When a significant wind is blowing over the surface, the mixing-length at the surface must increase to reflect larger surface oscillations than expected. Thus, a smaller value of δ is used, (Martin and Delhez, 1994).

Stratification inhibits vertical motions and turbulent exchanges. Therefore, to model such a process, the mixing length is reduced according to the intensity of the stratification, measured by the flux Richardson number (R_f), (Eq. 4.4, Table 1, 2). The values of ϵ and \tilde{v} are then obtained using simple dimensional arguments from the basic variables of the turbulence scheme (Eqs. 4.2 and 4.5, Table 1, 2). Turbulent diffusivities for scalar quantities are supposed to be proportional to the turbulent viscosity according to the Reynolds analogy (Eq. 4.6). These coefficients are also a function of R_f because the momentum exchange, involving pressure forces, is likely to maintain its efficiency at high Richardson numbers, whereas the transfer of the scalar quantities (temperature, T , Salinity, S and buoyancy, b) is reduced as stratification increases (Nihoul *et al.*, 1989).

Numerical aspects

The numerical model is based on a finite volume approach using the Arakawa C-grid (Arakawa and Lamb, 1977). For most of the terms in the equations,

TABLE 1. – Mathematical model of the general circulation.

Continuity equation

$$\nabla_H \mathbf{u} + \frac{\partial w}{\partial z} = 0 \quad (1)$$

Momentum equation

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla_H \mathbf{u} + \omega \frac{\partial \mathbf{u}}{\partial z} + f \mathbf{e}_3 \wedge \mathbf{u} = -\nabla_H q + \frac{\partial}{\partial z} \left(\tilde{\nu} \frac{\partial \mathbf{u}}{\partial z} \right) + K_H \nabla_H^2 \mathbf{u} \quad (2)$$

$$f = 2\Omega \sin(\text{lat}) \quad (2.1)$$

$$q = \frac{p}{\rho_0} + gz \quad (2.2)$$

$$\frac{\partial q}{\partial z} = b(T, S) = -g \frac{\rho - \rho_0}{\rho_0} \quad (2.3)$$

Equations for scalar constituents and properties – temperature and salinity:

$$\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla_H T + \frac{\partial}{\partial z} (wT) = \frac{\partial}{\partial z} \left(\tilde{\lambda} \frac{\partial T}{\partial z} \right) + K_H^S \nabla_H^2 T \quad (3.1)$$

$$\frac{\partial S}{\partial t} + \mathbf{u} \cdot \nabla_H S + \frac{\partial}{\partial z} (wS) = \frac{\partial}{\partial z} \left(\tilde{\lambda} \frac{\partial S}{\partial z} \right) + K_H^S \nabla_H^2 S \quad (3.2)$$

Turbulent closure

$$\frac{\partial k}{\partial t} + \mathbf{u} \cdot \nabla_H k + \frac{\partial}{\partial z} (wk) = Q^k + \frac{\partial}{\partial z} \left(\tilde{\lambda}^k \frac{\partial k}{\partial z} \right) \quad (4)$$

$$Q^k = \tilde{\nu} \left| \frac{\partial \mathbf{u}}{\partial z} \right|^2 - \tilde{\lambda}^b \frac{\partial b}{\partial z} - \varepsilon = \tilde{\nu} \left| \frac{\partial \mathbf{u}}{\partial z} \right|^2 (1 - R_f) - \varepsilon \quad (4.1)$$

$$\varepsilon = \frac{\alpha^2 k^{3/2}}{\ell} \quad (4.2)$$

$$\ell = \kappa(z+h) \left(1 - \delta \frac{z+h}{\xi+h} \right) \left(1 - \frac{R_f}{R_{fc}} \right) \quad (4.3)$$

$$R_f = \frac{\tilde{\lambda}^b \frac{\partial b}{\partial z}}{\tilde{\nu} \left| \frac{\partial \mathbf{u}}{\partial z} \right|^2} \quad (4.4)$$

$$\tilde{\nu} = \alpha \sqrt{k} \quad (4.5)$$

$$\tilde{\lambda}^T = \tilde{\lambda}^S = \tilde{\lambda}^b = \gamma \sqrt{1 - \frac{R_f}{R_{fc}}} \tilde{\nu}; \quad \tilde{\lambda}^k = \tilde{\nu} \quad (4.6)$$

the horizontal discretisation is centred in space. The advection terms of momentum, salinity and temperature are represented by means of a hybrid scheme in which the level of upwind is controlled locally according to the estimated curvature of the field (Delhez and Martin, 1992; James, 1986).

TABLE 2. – List of symbols used in Table 1.

t	Time (s)
z	Vertical coordinate above the reference level of the free-surface (m)
h	Unperturbed depth (m)
ξ	Free-surface elevation (m)
$H = h + \xi$	Total depth (m)
σ	Normalised vertical coordinate
\mathbf{v}	Three-dimensional velocity vector (m s ⁻¹)
\mathbf{u}	Horizontal current (m s ⁻¹)
∇_H	Horizontal differential operator
\mathbf{e}_3	Unit vertical vector pointing upwards
f	Latitude dependent Coriolis parameter (s ⁻¹)
K_H, K_H^S, K_H^T	Horizontal diffusion coefficients for momentum, salinity and temperature (m ² s ⁻¹)
g	Vertical acceleration of gravity (m s ⁻²)
ρ, ρ_0	Density of sea water and the reference value of the density (kg m ⁻³)
$b(T, S)$	Buoyancy (m s ⁻²)
T	Temperature (°C)
S	Salinity
k	Turbulent kinetic energy per unit mass (associated with all sub-mesoscale processes) (m ² s ⁻²)
ε	Dissipation rate of k (m ² s ⁻³)
Q	Macroscale production-destruction of turbulence
ℓ	Mixing length (m)
R_f	Macroscale Richardson flux number
$\tilde{\nu}$	Turbulent viscosity (m ² s ⁻¹)
$\tilde{\lambda}^b$	Turbulent diffusivity of buoyancy (m ² s ⁻¹)
$\tilde{\lambda}^T, \tilde{\lambda}^S, \tilde{\lambda}^k$	Turbulent diffusivity of temperature, salinity and turbulent kinetic energy (m ² s ⁻¹)
$\alpha, \delta, \gamma, \kappa, R_{fc}$	Model parameters

The time stepping uses the mode splitting technique (e.g. Delhez and Martin, 1992) according to which the depth-integrated continuity and momentum equations are first solved with a small time step, allowing for a good representation of the dynamics of the fast-moving surface gravity waves, and then a larger time step is used to solve the full set of 3D equations in which the barotropic component of the pressure gradient (i.e. the sea surface slope) is taken from the solution of the depth-integrated mode. For the depth-integrated equations, the external mode, an ADI scheme is used (Beckers and Neves, 1985). For the 3D equations, the internal mode, the discretization is explicit except for the vertical advection and diffusion terms that are taken implicitly for stability reasons (Beckers and Deleersnijder, 1993). As a result, a time step of 15 s is used to solve the external node without any significant phase error while the (more demanding) internal dynamics is addressed with a time step of 150 s according to the CFL condition.

The code also includes a new “drying and wetting” algorithm to cope with the possible drying of sandbanks at low tide. This algorithm is implemented in the solution of the external model. A grid cell is simply removed from the computational domain when the water height decreases below a given threshold (H_{dry} ,

TABLE 3. – Values of the hydrodynamic parameters in the numerical model

Hydrodynamic parameter	Numerical value
Time step. T2D, T3D.	15 s, 150 s
Horizontal viscosity.	100 m ² s ⁻¹
Horizontal diffusivity of temperature and salinity.	10 m ² s ⁻¹
Horizontal diffusivity of turbulent kinetic energy	100 m ² s ⁻¹
<i>Turbulent closure parameters</i>	
α	0.5
γ	1.1
κ (Von Karman constant)	0.4
δ (strong wind conditions)	0.5
Roughness length z_0	4.19E-3 m
Minimum drag coefficient. C_{DMIN}	0.0025
H_{dry} , H_{flood}	0.05 m, 0.05 m

Table 3) and re-introduced later when the sea surface slope between the dried cell and the surrounding cells exceeds another threshold (H_{flood} , Table 3).

Horizontal and vertical discretisation of the domain

The numerical grid extends from 42.14°N to 42.25°N and from 8.92°E to 8.65°E (Fig. 1). It is oriented along the NE-SW axis of the Ría in order to provide a good representation of the large-scale bathymetric structure and hence of the flow in and out. Steeper horizontal variations are expected to develop more across the Ría than along the axis, and consequently the mesh is regular in the horizontal but has a grid size of 300 m in the direction across the Ría vs. 500 metres along the axis. This yields 42 x 53 grid boxes per horizontal level. Notice, however, that only about half of them are actually wet points.

The vertical discretization is done in the transformed σ -coordinate system where σ , the new vertical coordinate, is defined as:

$$\sigma = \frac{z+h}{\xi+h} \in [0,1]$$

This transformation allows a continuous description of the bathymetry and hence a good representation of this important forcing factor, particularly for shallow water motions. Ten vertical levels are used in the σ space. They are unevenly spaced to provide a finer resolution in which large vertical variations can be anticipated, i.e. at the surface and at the level of the possible thermocline. The distribution of $\Delta\sigma$ is, from bottom to top:

0.05/0.1/0.1/0.15/0.15/0.15/0.15/0.1/0.1/0.1/0.05/0.05

The real thickness ($\Delta z = H \Delta\sigma$) of the different levels adapts to the local depth so that the discretization in the σ coordinate system is valid for the whole domain. The thickness of the first level $\Delta\sigma$ at the bottom is taken small enough for the first U- and V-points chosen to be in the logarithmic bottom boundary layer in the vast majority of the domain. The effect of the bottom boundary layer can then be easily and economically parameterized.

Boundary conditions

Surface boundary conditions

Boundary conditions at the air-sea interface are obtained by expressing those turbulent vertical fluxes of momentum, heat, etc. that are continuous except for boundary sources or sinks.

Turbulent Kinetic energy flux

The flux of turbulent kinetic energy at the surface boundary depends on wind intensity (Eq. 2, Table 4). Without wind, there is no input of k through the surface. On the other hand, when the wind is blowing over the sea, breaking waves inject turbulence into the water column. We therefore preferred to prescribe the flux of turbulent kinetic energy instead of the surface value of k , because this allows turbulent kinetic energy in the surface boundary layer to depend not only on wind conditions but also on what happens below.

Momentum flux

As the turbulent flux, the momentum flux in the atmospheric boundary layer depends on the wind forcing (Eq. 1, Table 4). Two different wind fields are used inside the domain to reproduce the differences in wind intensity and direction that are found inside and outside the Ría. The outer zone (from the open boundary to Home Cape, Fig. 1) is under a much greater oceanic influence and the winds blowing in this region are more intense, with a marked trend that could induce a long-term transport. In the innermost zone (from Mar Cape to the strait of Rande), the wind field is weaker and generally blowing along the main axis following the Ría's orography. Between the two zones, the model interpolates the two wind fields linearly.

TABLE 4. – Hydrodynamic boundary conditions.

Momentum surface flux

$$\tilde{\nu} \frac{\partial \mathbf{u}}{\partial z} = C_0 (1 + 0.1 |\mathbf{W}|) |\mathbf{W}| \mathbf{W} \quad (1)$$

Turbulent kinetic energy surface flux

$$\tilde{\nu} \frac{\partial k}{\partial z} = 3C_0 10^{-3} |\mathbf{W}|^3 \quad (2)$$

$$C_0 = 0.63 \cdot 10^{-6}$$

$\mathbf{W} = (u_0, v_0)$ surface wind velocity in m s^{-1}

Heat surface flux

$$Q_0 = \rho c_p \lambda \frac{\partial T}{\partial z} = I + C - IR - E \quad (3)$$

$$I = [3 \cdot 191 + 1 \cdot 115 \sin^2 [\pi(355 - J)/365]] [50 \cdot 417 - 4 \cdot 474 N]$$

$$IR = (143 - 0.9 T_s - 0.46 e_a)(1 - 0.1 N)$$

$$C = 24.88(0.38 + 0.114 |\mathbf{W}|)(T_s - T_a)$$

$$QE = 58.7 E$$

$$E = (0.26 + 0.077 |\mathbf{W}|) (e_s - e_z)$$

ρ	Density of the water
C_p	Specific heat of the water
Q_0	Surface heat flux
I	Solar radiation (w m^{-2})
J	Day number starting 1 January
N	Cloud cover in oktas
IR	Infrared back radiation (w m^{-2})
T_s	Water surface temperature ($^{\circ}\text{C}$)
H_a	Relative humidity above the water surface
e_a	Vapour pressure 2 m above the water surface (mmHg)
e_s	Vapour pressure of water at the sea surface (mmHg)
C	Heat by conduction
T_a	Air temperature
QE	Heat lost by evaporation (w m^{-2})
E	Evaporation (mm day^{-1})

Salinity surface flux

There is no salinity flux at surface.

Freshwater open boundary discharge

$$Q = 0.007538 + 0.007889 A \quad (4)$$

$$A = (1 - k) / \left(k - k^{365+1} \right) \sum_{n=1}^{365} P_n k^n$$

Q	Flux of fresh water in $\text{m}^3 \text{s}^{-1} \text{km}^{-1}$
A	Drainage expressed in $\text{l m}^{-2} \text{d}^{-1}$
P_n	Rainfall collected during day n, in l m^{-1}
k	Retention coefficient en $\text{m}^{-1} \text{d}^{-1}$

Momentum bottom flux

$$|\tau_b| = C_D |\mathbf{u}_b| \mathbf{u}_b \quad (5)$$

$$C_D = \left(\left[\frac{\kappa}{\ln(z/z_0)} \right]^2, 0.0025 \right)$$

z_0	Roughness length.
κ	0.4 Von Karman constant.
C_D	Bottom drag coefficient
\mathbf{u}_b	Bottom current velocity in m s^{-1}

Heat and salinity bottom fluxes

There are no heat and salinity fluxes at the bottom.

Heat Flux

The exchange of heat between the sea surface and the atmosphere is calculated by a set of well-known equations (i.e. Rosón *et al.*, 1997), taking into account solar radiation, back radiation, conduction and evaporation (Eq. 3, Table 4). Incoming solar radiation is estimated using Mosby's formula (i.e. Dietrich *et al.*, 1980), and the infrared back radiation component and the exchange of heat by conduction and heat lost by evaporation are evaluated by empirical equations (Otto, 1975). Air temperature, relative humidity and cloud cover are then required to compute the heat flux through the sur-

face.

Salinity

We assume that the precipitation and evaporation balance each other. The salt flux through the surface is thus supposed to be zero.

*Open boundary conditions**2D Elevation boundary conditions*

At the open boundary with the shelf, M_2 , S_2 , N_2 and L_2 , the most important tidal constituents for the

Ría of Vigo are introduced through a radiation condition based on the theory of characteristics (Hedstrom, 1979; Røed and Cooper, 1987) and applied to the depth-integrated part of the signal. This procedure assumes that one Riemann variable (a combination of the total transport and free-surface elevation) is prescribed from information taken outside the domain, while the other Riemann variable is radiated outwards. The actual boundary value arises from the recombination of the two Riemann variables (Røed and Cooper, 1987). This procedure allows the forcing data from outside the domain to be incorporated, whilst avoiding spurious reflection on the boundary of perturbations generated inside the Ría.

Temperature and salinity

Salinity and temperature of inflowing shelf water are prescribed using observed data (real observations or climatological mean). At outflow boundaries, where water leaves the domain, the gradients of salinity and temperature normal to the boundary are supposed to vanish.

Freshwater discharge

The input of fresh water coming out of the Oitabén and Verdugo rivers is introduced by imposing this inflow on the top 5 meters of the water column at the inner open boundary. This freshwater runoff is estimated as a function of precipitation by using the equation given by Ríos *et al.* (1992) (Eq. 4, Table 4).

Solid boundary conditions

The coastline is a streamline and we must thus fix its impermeability. The component of the velocity vector normal to the coastline is set to zero at all depths. By virtue of their impermeability there are no advective or diffusive fluxes of any variable through solid boundaries.

As the bottom layer has been chosen to be in the logarithmic bottom boundary layer, the bottom stress can be parameterised using a quadratic expression of the bottom velocity, (Blumberg and Mellor, 1987) (Eq. 5 in Table 4). In those places where the bottom boundary layer is not well resolved (deepest areas), the drag coefficient is set to a minimum value $C_D=0.0025$ (c.f. Martín *et al.*, 1993). A lateral friction is further introduced to take into account the general slowing down of the stream induced by the unresolved coastal

boundary layer.

SIMULATION CONDITIONS

The model was run to simulate the typical conditions that can be found at the end of the summer when changes between North and South winds are often taking place, defining two markedly different scenarios:

a) normal (or normal reinforced) estuarine circulation, with low salinity surface water flowing outwards and denser NACW water coming inwards in the bottom layer;

b) negative circulation due to intense south winds, with marked thermohaline fronts and near vertically homogeneous water conditions at the Ría's mouth.

In both situations, a constant river discharge of $13 \text{ m}^3 \text{ s}^{-1}$ was assumed (the last ten years' mean value for the first fortnight of September). Rainfall used to calculate the river runoff was monitored by the Instituto Nacional de Meteorología at the meteorological station in Peinador airport, and corrected taking into account its altitude above sea level, (Ríos *et al.*, 1992). Air temperature, relative humidity and cloud cover were provided by a meteorological station situated close to the Ría (Fig. 1).

The model started from rest with initial temperature and salinity (16°C , 35.3 and 14°C , 35.7 respectively) based on the observed hydrographic data taken at the isobath of 60 m during the September 1998 INCOCEANO 2 survey.

Temperature and salinity open boundary conditions were related to the observed hydrographic profiles measured at the South mouth of the Ría (Fig. 1) in different wind conditions during the INCOCEANO 2 survey. The same boundary conditions were imposed at both the southern and the northern mouth.

Before the results were collected, a 30 day spin-up was used under near negligible northerly wind to guarantee the dynamic adjustment of the flow field and the thermohaline structure. The resulting fields are characterised by a weak estuarine type circulation pattern with a residual current of about 0.04 m s^{-1} .

In order to appraise the influence of different meteorological conditions on the hydrodynamics, two numerical experiments were then carried out using two types of wind forcing: (1) upwelling-favourable winds and (2) downwelling-favourable winds.

In the outer zone, the upwelling-favourable condition referred to a wind blowing southward (350°), while the downwelling-favourable situation brought

up the wind blowing near northward (190°). The magnitude of the wind was set to 7 m s^{-1} in both cases, which correspond approximately to the mean of intense wind periods taking place during August and September. In the inner zone, the wind speed is taken to be half of the value on the shelf, while the direction reflects the channelling effect of the Ría with southwestward winds (40°) in an upwelling-favourable situation and northeastward winds (220°) in an downwelling-favourable situation.

Both simulations have been performed with the same fluctuations in the solar radiation, cloud cover, relative humidity and fresh water inputs, which are representative of the natural variability observed in the data sets during the first fortnight of September.

RESULTS AND DISCUSSION

Upwelling-favourable wind scenario

Current field

After a period of 5 days with upwelling favourable winds (a typical periodicity for the wind variations in September (Figueiras *et al.*, 1994), the current field, and accordingly the temperature and salinity fields, were significantly altered. The low salinity water was pushed away out of the Ría by the offshore surface Ekman-transport. The deep saltier and denser waters flowed into the Ría along the bottom. The result was a general pattern of residual circulation in two layers that is typical of a positive estuarine circulation and has been described previously by several authors (Mouriño *et al.*, 1984; Blanton *et al.*, 1984; Prego and Fraga, 1992) (Fig. 2).

At the surface, the wind induced an out-going residual current that fluctuated around $0.13 \pm 0.04 \text{ m s}^{-1}$ in the surface layer and progressively diminished to reach zero at the level of the thermocline, (Fig. 3f). Taboada *et al.* (1998) gave smaller outgoing velocities, probably due to the lack of an explicit wind forcing mechanism in their model. In fact, surface velocities estimated from real hydrological data and taking into account the wind forcing showed an average surface velocity of 0.13 m s^{-1} during a September upwelling event (Alvarez-Salgado *et al.*, 1998). Also, a near-surface current meter record deployed at 2.7 m depth from 11 to 27 September 1990 in the middle of the Ría de Vigo (Figueiras *et al.*, 1994) showed a persistent outflow of water (average of 0.09 m s^{-1}) to the southwest dur-

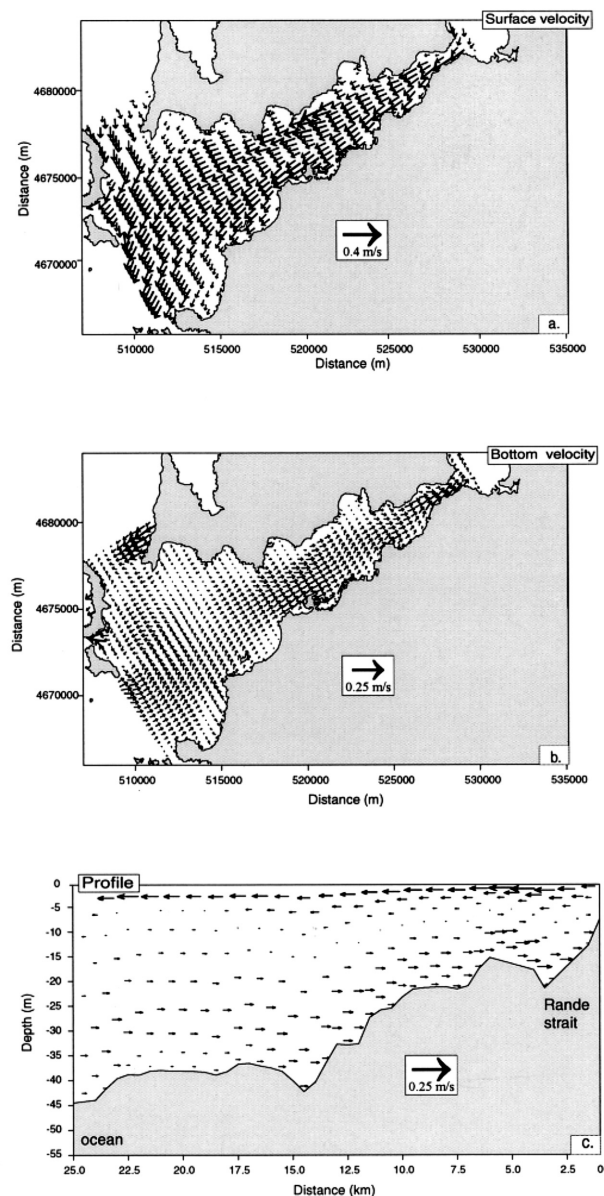


FIG. 2. – Residual circulation during upwelling-favourable wind (a) in the surface layer and (b) in the bottom layer. (c) Vertical profile of the residual velocity component u plotted along the main axis of the Ría.

ing an upwelling situation.

The most striking feature of the surface circulation was the moderately strong (around 0.15 m s^{-1}) current leaving the Ría at its Southwest mouth, in contrast with a weak (0.03 m s^{-1}) stream entering through the northern mouth. Neither the box model of Prego *et al.* (1992) nor the 3-D hydrodynamical model of Taboada *et al.* (1998) reproduced this peculiarity. The bathymetry and differences in orientation of the two mouths with respect to wind direction could influence in the general circulation: when quite strong northerly winds are blowing, the

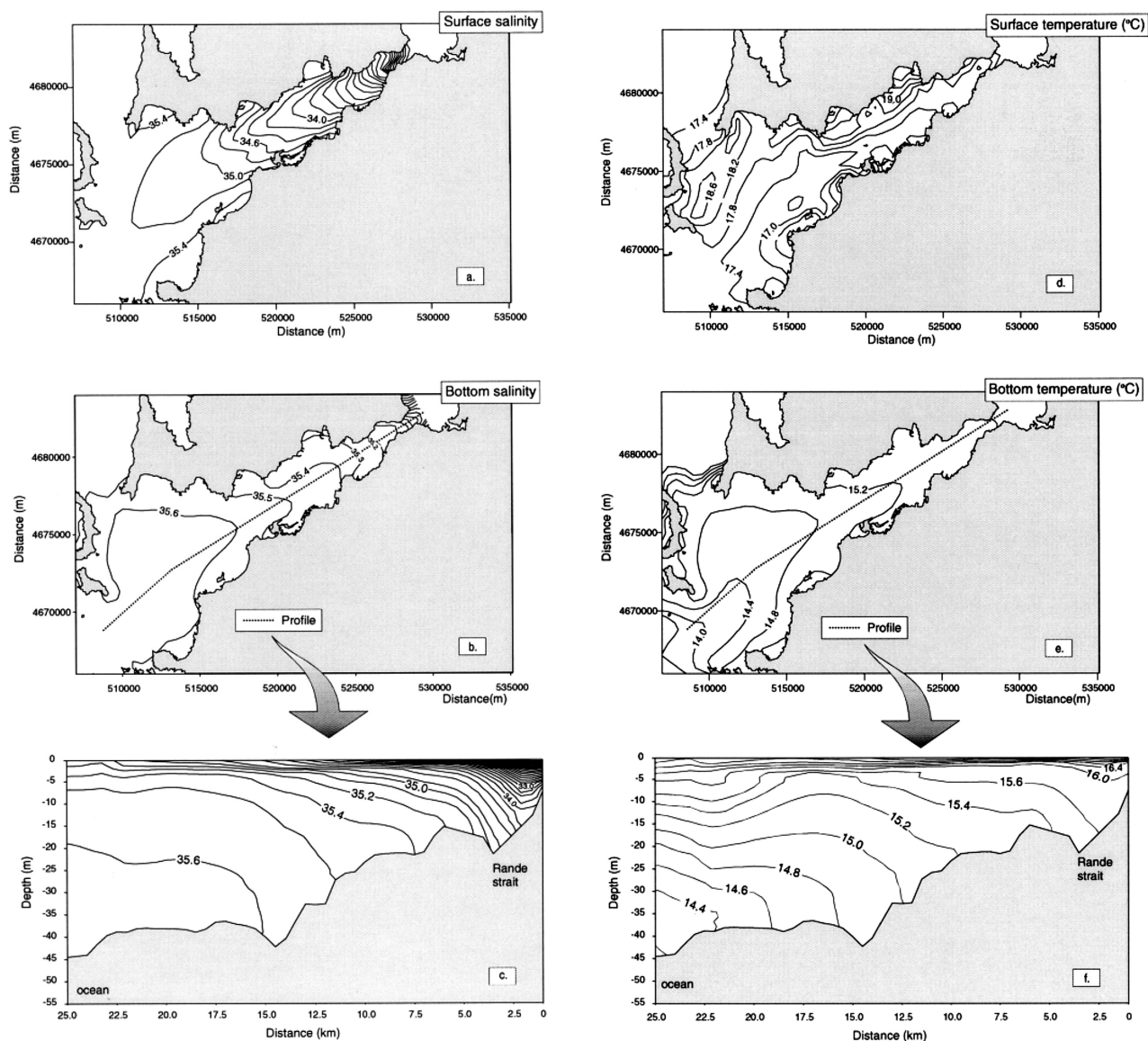


FIG. 3. – Salinity and temperature distribution after 5 days of a moderately strong (7 m s^{-1}) northerly wind. (a) surface salinity, (b) bottom salinity, (c) vertical salinity profile depicted in the dotted line in (b), (d) surface temperature, (e) bottom temperature, and (f) vertical salinity profile as shown in the dotted line in (e).

southwestern mouth (which is 5 km wide and 67 m deep) forms a free open boundary, where a clear outgoing signal could be developed by means of the “wind-induced” transport. On the other hand, in the smaller northern mouth (only 2.5 km wide and 23 m deep) the wind-induced surface transport coming from the continental shelf blocks out the water flowing outwards and causes it to turn to the wider and deeper southern mouth. Nevertheless, some unrealistic values are detected at the northern mouth, probably caused by still unresolved boundary problems.

Below the thermocline, the wind forcing had nearly no influence on the general circulation and a compensatory inshore flux was established below

the zero velocity level, as is shown in Fig. 2b. The incoming saltier water penetrated through the southwestern mouth and forked into two different currents: one flowing in the direction of the main axis of the Ría, and one turning to the right and describing a cyclonic gyro (already cited by Fraga and Prego, 1989) and probably induced by the presence of the Cíes Islands.

In the bottom layer, the magnitude of the velocity ranged from 0.01 m s^{-1} (where there was no influence of the incoming current) to 0.05 m s^{-1} in the rest of the domain. Both the bottom circulation pattern and the average velocities were in close agreement with those given by the model of Taboada *et al.*

(1998). Some observational data support both model results, since Álvarez-Salgado *et al.* (1998) found depth-averaged residual velocities (4–26 m) of 0.045 m s^{-1} at the middle of the Ría's main axis.

Over the studied period, the transition zone between opposite flows (level of non-horizontal motion) reflected the bathymetry (Fig. 2c). It uplifted inward from 8.5 m at the open boundary to 4.5 m at the inner part, getting deeper in the middle (14 m). This near-zero velocity level closely follows the thermocline depth (Fig. 3c) and agrees reasonably with the one simulated by the model of Taboada *et al.* (1998) (7.6 m in the inner part to 8.9 m in the outer zone) or that of Prego *et al.* (1992) (5.4 m in the inner box and 9.3 m in the outer box).

Moreover, Álvarez-Salgado *et al.* (1998), and Prego and Fraga (1992) found similar results in September observed data. Both of them localised the non-motion level (assuming it as the pycnocline depth) at a depth ranging from 4 m to 6 m during different upwelling events.

To sum up, a north wind episode results in a single circulation cell with a weak upwelling in the middle part of the Ría limited by the strong stratification, a narrow and strong offshore flow at the surface, and a compensatory wider and weaker onshore flow in the deep region.

Temperature and salinity fields

There were appreciable quantitative differences in the salinity and temperature surface layer between one side of the Ría and the other. Low salinity values (<32) were observed in the vicinity of San Simón bay, where rivers discharge. This fresh water flows out preferably along the north shore. Warm water ($>18^\circ\text{C}$) followed this low saline stream along the north shore until Home Cape and then turned to the southern mouth. On the other side, at the south shore, saltier and colder water (>35.4 , $<16.60^\circ\text{C}$) uplifted from deeper levels and rose to the surface (Figs. 3a and 3d).

In the bottom layer, lower salinities (<35.0) and higher temperatures ($>15.40^\circ\text{C}$) were located close to the Rande Strait, while saltier and colder waters (>35.6 , $<14.0^\circ\text{C}$) penetrated through the southern mouth and flowed into the Ría close to the south shore (Figs. 3b and 3e).

The upwelling situation was evident in the thermal structure through the uplift of the isotherms and the compression of the temperature gradient in the upper 5 m (Fig. 3f). Although relatively high temperatures

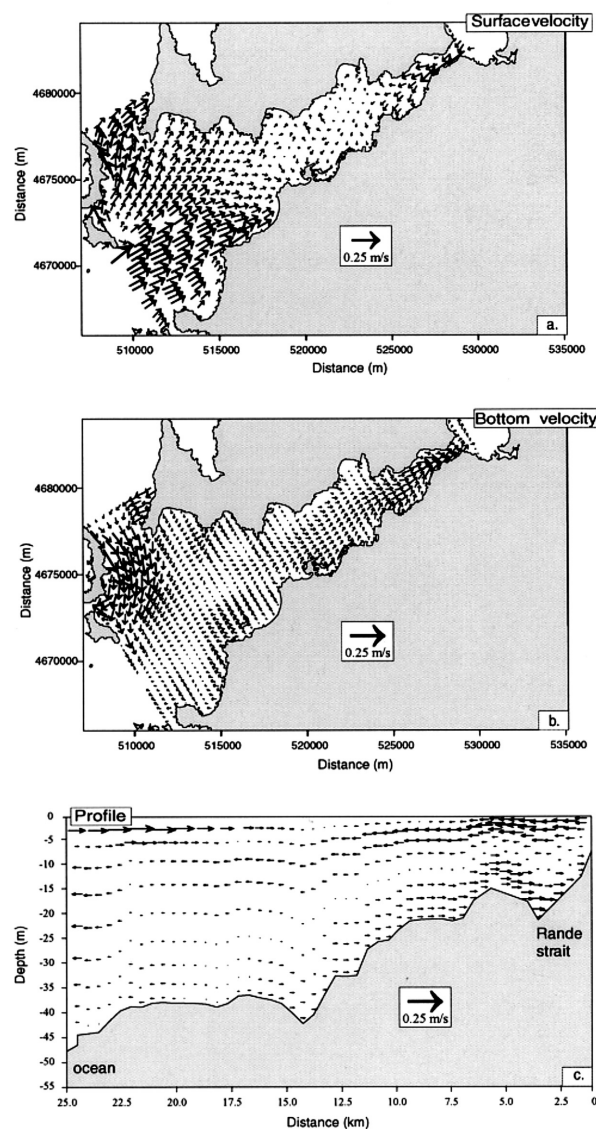


FIG. 4. – Residual circulation during downwelling-favourable wind (a) in the surface layer and (b) in the bottom layer. (c) Vertical profile of the residual velocity component u plotted along the main axis of the Ría.

($>14^\circ\text{C}$) have been measured in the bottom layer during upwelling events and were justified by the rapid mixing with other water bodies when cold shelf water enters the Ría (Álvarez-Salgado, 1996), it is more usual to find colder waters ($<13.5^\circ\text{C}$). The predicted bottom temperatures, which ranged from 14.4°C to 15.5°C (Fig. 3f), were slightly higher than those expected from observations. The differences may be explained by the fact that the domain did not include a greater percentage of the adjacent shelf and the entrance of shelf cold water caused by an upwelling event off the Galician coast was possible underestimated. Nevertheless, the temperature and salinity model results compared reasonably well (especially in the upper layer) with numerous observational data

(e.g. Figueiras *et al.*, 1994; Egleé, 1996).

Therefore, the combined effect of temperature and salinity resulted in the development of a strong stratification that led to a positive circulation. The thermocline and the halocline depth coincided with the non-motion level, ranging from 3.5 m in the inner part to 8 m in the outer part.

Downwelling-favourable winds scenario

Current field

During a downwelling situation, south winds advected relatively dense oceanic surface water into the Ría, generating a major onshore transport. Con-

sequently, surface waters from the inner part of the Ría were strongly mixed with shelf ones, leading to a convergence front 20 km off the Rande Strait, as isotherm and isohaline deepening indicates (Figs. 4 and 5), and dividing the Ría into two different circulation cells with different thermohaline properties.

In the inner and middle part of the domain, the spatial structure of the residual circulation was similar to that generated by northerly winds. This area maintained a positive estuarine circulation, although the magnitude of the velocities had diminished strongly due to the blocking of the surface outflow at the front zone; the residual velocity was less than 0.03 m s^{-1} throughout the column, increasing the water flushing time (9.5 days) inside the Ría. The

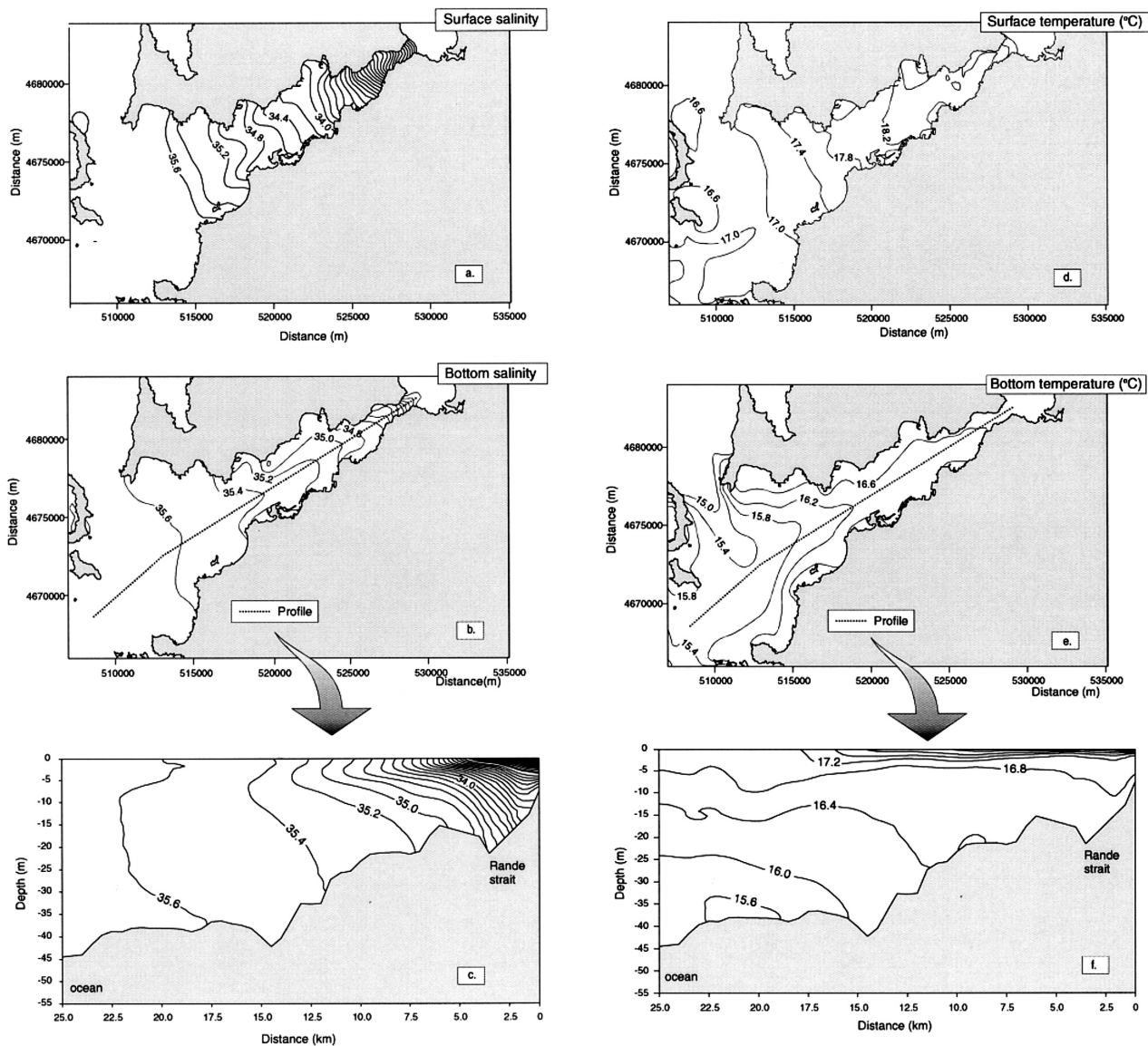


FIG. 5. – Salinity and temperature distribution after 5 days of a moderately strong (7 m s^{-1}) southerly wind. (a) surface salinity, (b) bottom salinity, (c) vertical salinity profile depicted in the dotted line in (b), (d) surface temperature, (e) bottom temperature, and (f) vertical salinity profile as shown in the dotted line in (e).

position of the non-motion level in this zone was deeper than the zero velocity level generated by North winds, uplifting inward from 30 m at the frontal zone to about 10 m close to San Simon Bay. The result was the development of a semi-closed circulation pattern (Fig. 4c).

Nearer to the shelf, the associated residual currents form a characteristic inverse circulation pattern, with a strong water inflow at the surface through the southern mouth ($0.14 \pm 0.02 \text{ m s}^{-1}$) and a weaker and wider countercurrent ($0.02 \pm 0.02 \text{ m s}^{-1}$) between 10 m and 30 m depth. A similar circulation pattern has been reported with South wind episodes (Figueiras *et al.*, 1994; Tilstone *et al.*, 1994; Fraga and Prego, 1989; Castro *et al.*, 1994).

Few real and no modelled velocity data were available to compare model results in a downwelling situation. Using a surface current meter, Figueiras *et al.* (1994) recorded a reversed net residual flow into the Ría to the Northeast, in accordance with the residual velocity fields of the model. The magnitude of the velocity ranged between 0.11 m s^{-1} and 0.01 m s^{-1} (with a mean of 0.05 m s^{-1}).

Temperature and salinity field

The downwelling-favourable thermohaline distribution differed markedly from the upwelling-favourable one. South wind enhanced the turbulent vertical mixing and led to a nearly uniform vertical structure of thermohaline properties.

In the surface layer (Figs. 5a and 5d), the lowest salinity and highest temperature values (32 and 19°C) corresponding to bodies of less saline water due to river runoff were found in the innermost area, where river water was kept back by the shelf influx. The saltiest and coldest waters (35.6 and 17.0°C) were clearly present at the most external part of the Ría.

At the bottom, the salinity field remained quite invariable in comparison with the surface one, and somewhat saltier (33.0) at the head of the Ría (Fig. 5c). Indeed, the halocline was not so developed as with north winds (Fig. 3c).

Changes in the vertical temperature and salinity distribution in different wind conditions were evident and in close accordance with observations (Figueiras *et al.*, 1994). During upwelling-favourable winds the isotherm uplift and warm water was confined near the surface. North wind tended to speed up the interchange with the shelf, diminishing the water flushing time inside the Ría (4.5 days) and therefore reducing the warming of the

upper layer. In contrast, during downwelling-favourable winds the colder water observed in the initial conditions moved out towards the ocean, leaving a more homogeneous water column with temperatures of $>15.8^\circ\text{C}$. Moreover, a general warming took place due the increasing of the water flushing time (9.5 days) inside the Ría.

CONCLUSIONS

The use of a mathematical 3-D model to investigate the hydrodynamic processes in the Ría de Vigo provides a considerable help in the description of observations. Observed features can be explained from the point of view of their dynamics and their links with various processes and forcing factors.

This model pushes forwards the study of the hydrodynamic structure in the Ría de Vigo with different meteorological forcing, inducing noticeable changes in the thermohaline distribution and the residual circulation pattern at a time scale of the order of five days. This is an important advance in comparison with earlier model attempts (e.g. Taboada).

A downwelling situation has never been simulated using a 3-D model. Winds with a strong north component tend to produce a positive estuarine circulation, while powerful southerly winds arise in an atypical reversal circulation pattern, which reduces the outgoing velocities inside the Ría. These results are supported by the fact that salinity, temperature and velocity fields are in good agreement with the experimental data.

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