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**Decomposing the intra-annual variability of flushing characteristics in a tidal bay
along the North Sea**

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Highlights:

- Large spatial and intra-annual variability in turnover times of the Oosterschelde
- Winds impact the local turnover time by up to ± 50 days
- Density gradient and turnover time are regulated by neighboring river outflows

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Abstract

Flushing timescales in estuaries and coastal bays largely shape the distribution of dissolved and particulate matter and therefore have important environmental implications. Here we investigate the spatiotemporal variability of turnover times in a semi-enclosed tidal bay and examine the potential underlying causes. The basin considered here is the Dutch coastal bay Oosterschelde (Eastern Scheldt in English) on the east coast of the North Sea. Using a calibrated and validated 3D hydrodynamic model, conservative tracer experiments were conducted to estimate turnover times. Sensitivity scenarios were run to examine the relative roles of tides, winds, and gravitational flow in driving the intra-annual variability. Results indicate that gravitational flow and winds contribute most to the intra-annual variability of turnover times in the Oosterschelde. Freshets from the Westerschelde and Rhine Rivers extending to the mouth of the bay can reduce the axial density gradient and gravitational flow, increasing turnover times of the basin. The prevailing southwesterly winds, when of sufficient magnitude (> 5 m/s), modify the spatial patterns of turnover times especially in shallow waters. Tides force flushing of the bay but do not induce significant intra-annual variability. Findings in this study underline the importance of intra-annual variability of turnover times in coastal bays and methodology proposed here facilitates intercomparison studies for other systems.

Keywords

Turnover time; tides; winds; gravitational force; intra-annual variability; tidal bay

1 Introduction

Estuaries and coastal bays are subject to natural and anthropogenic influences (e.g., sea level rise, eutrophication, pollution), posing growing threats to the sustainability of these systems. The spatiotemporal distribution of environmentally important dissolved and particulate substances (e.g., nutrients, chlorophyll-a, dissolved oxygen, trace metals, bacteria) can be interpreted by considering flushing timescales such as residence time, water age, turnover time, flushing time, renewal time, etc. (Banas et al., 2015; Guyondet et al., 2005; Josefson & Rasmussen, 2000; Nixon et al., 1996; Prandle, 1984; Shen & Wang, 2007; Soetaert & Herman, 1995; Torr  ton et al., 2007; Wan et al., 2013). For example, a system with a flushing rate faster than the phytoplankton growth rate is relatively resistant to eutrophication and harmful algal blooms (Howarth et al., 2000; Wang et al., 2004). In ecosystems with intensive bivalve aquaculture, the production carrying capacity is considered to be exceeded when the residence time is longer than the clearance time needed for the cultured shellfish to filter the entire water column (Gibbs, 2007; Koutitonsky et al., 2004). In addition, system-wide flushing timescales are also applied as quantitative indicators to intercompare and categorize estuaries and coastal lagoons, as well as to divide basins into sub-regions (Abdelrhman, 2005; Ferrarin et al., 2013; Lemagie & Lerczak, 2015; Plus et al., 2009; Umgieser et al., 2014). Therefore, estimating flushing timescales is not only an essential step to understand transport processes in nearshore systems, but also of significant biogeochemical and ecological relevance.

Flushing of estuaries and coastal bays is a complicated process forced by freshwater runoff, tides, meteorological conditions, gravitational circulation, and their interaction with bathymetry and geometry of the system (B  rcena et al., 2012; Defne & Ganju, 2015). River discharge is usually the first-order contributor to transport and flushing in many freshwater-fed estuaries (Du et al., 2018; Lemagie & Lerczak, 2015; Shen & Wang, 2007), while in tidal bays or lagoons where river runoff is limited, tides are the major force driving the water renewal (Cucco et al., 2009; Viero & Defina, 2016; Yuan et al., 2007). Winds and

gravitational forces have the potential to modulate transport processes via complex mechanisms (Banas & Hickey, 2005; Choi & Lee, 2004; Cucco & Umgiesser, 2006; Tartinville et al., 1997; Webb & Marr, 2016). Furthermore, the shape, geometry, and bathymetry of estuaries/lagoons can add to the spatial heterogeneity of flushing characteristics (Banas & Hickey, 2005; Ferrarin et al., 2013; Umgiesser et al., 2014). Therefore, a systematic study of various processes is required to understand the flushing dynamics in these systems.

Dating back to Bolin & Rodhe (1973), concepts like age, transit time, turnover time, and residence time have been used to quantify different kinds of renewal timescales in a reservoir. However, in theoretical and realistic oceanographic studies, they have been adopted with varying definitions and calculation methods. Lemagie & Lerczak (2015), Monsen et al. (2002), Rayson et al. (2016), Viero & Defina (2016), and many other studies reviewed and compared the main flushing timescales with practical examples and suggested that suitable timescales should be selected based on their definitions, computational costs, and the research objectives. Given the spatial heterogeneity in many systems, the bulk flushing timescales that treat the entire water body as a homogenous system (e.g., Lemagie & Lerczak, 2015) are not generally suitable, but rather spatially-variable (or local) transport timescales should be estimated with a 2D or 3D hydrodynamic model. To this end, both Eulerian and Lagrangian approaches are applicable (Cucco et al., 2009). However, due to the chaotic stirring in tide-dominated basins (Cucco et al., 2009; Ridderinkhof & Zimmerman, 1992), accurate estimates of transport timescales with Lagrangian particles require releases at various tidal phases, statistical averaging based on numerous particles at any control region, and high-order (random-walk) schemes to interpret diffusion (Luff & Pohlmann, 1995; Monsen et al., 2002; Oliveira & Baptista, 1997; Sandery & Kämpf, 2007). Thus, running intra-annual scenarios and sensitivity experiments with the Lagrangian method is computationally expensive (Rayson et al., 2016). Here we adopt turnover times, which represent the timescale of flushing a fraction (e^{-1}) of original water mass out of a basin, as defined by Zimmerman (1976) and further described in section 3.2.

In a way, turnover times are a characteristic of a basin or bay, but in this study we put forward the notion that they are not a constant, but, in fact, vary considerably both spatially and temporally. We focus on intra-annual variability and the underlying mechanisms. As an example, we investigate the intra-annual variability of flushing capability in the Oosterschelde (Eastern Scheldt in English), a semi-enclosed tidal bay in the southwestern Netherlands (Fig. 1) with limited freshwater input due to the construction of dams and sluices, which isolate it from the surrounding rivers (Nienhuis & Smaal, 1994). Given the computational efficiency and advantages in resolving diffusion, the Eulerian tracer approach is applied, using an existing 3D Oosterschelde hydrodynamic model driven by realistic external forcing (Jiang et al., 2019a). Using turnover time as an indicator, the magnitude and spatial patterns of intra-annual flushing variability in the Oosterschelde are delineated. Sensitivity experiments of major forcing agents were conducted to understand the contributors to the intra-annual variability. The study thus aims to explore the driving mechanisms behind variability of turnover times and to provide reference for biogeochemical and ecological studies.

2 The study site

The 350-km² Oosterschelde is located in the southwestern delta of the Netherlands and connected to the North Sea through a storm surge barrier at its mouth (Fig. 1). Of the three openings of the storm surge barrier (Fig. 1), which are only closed during extreme weather conditions (roughly once a year, but not closed during our study period 2009–2010), the southern opening and channel account for the most (60%) tidal flux into the basin (Eelkema et al., 2012). Vast shallow regions surrounding the tidal channels include a total of 110-km²

intertidal flats that are important wildlife habitats (salt marshes, seagrass, water birds, and bivalves, Tangelder et al., 2012). However, due to reduced tidal currents in the post-barrier period, intertidal flats have slowly but continuously eroded, and sediments are deposited in tidal gullies (Eelkema et al., 2012; Nienhuis & Smaal, 1994).

Prior physical and ecological studies usually divide the basin into four compartments (e.g., Jiang et al., 2019a; Smaal et al., 2001; Vroon, 1994; Wetsteyn & Kromkamp, 1994), the western, central, eastern, and northern (Fig. 1). We added a “mouth” region outside the storm surge barrier to characterize water masses entering the basin, where salinity is influenced by the outflow of the Westerschelde and Rhine rivers, as well as the North Sea water (Fig. 1). Tides in the Oosterschelde are mostly semi-diurnal. The mean tidal range calculated from both tidal measurements by the Dutch government agency Rijkswaterstaat (<https://waterinfo.rws.nl>) and modeled water elevation (Jiang et al., 2019a) indicates an increase from 2.5 m at the mouth to 3.4 m at the eastern end. Salinity and temperature temporal variations are largely controlled by the tidal motion flushing the North Sea waters in and out, and vertical stratification is very limited (Wetsteyn & Kromkamp, 1994). According to our recent cruise observations, the vertical salinity gradient is less than 0.017 m^{-1} . The maximum tidal velocity is around $1\text{--}1.5 \text{ m s}^{-1}$ in tidal channels (Mulder & Louters, 1994). Due to the dominant role of tidal flushing, there is a seaward increasing gradient of many substances along the basin, for example, salinity, Chl-a, seston, and turbidity (Nienhuis & Smaal, 1994; Wetsteyn & Kromkamp, 1994). The overall freshwater inflow into the bay is below $10 \text{ m}^3 \text{ s}^{-1}$ (Ysebaert et al., 2016). It was $3.2 \text{ m}^3 \text{ s}^{-1}$ and $4.5 \text{ m}^3 \text{ s}^{-1}$ in 2009 and 2010, respectively (Rijkswaterstaat data). This is a negligible amount compared to the flushing of the basin by tidal exchange, which is $\sim 2 \times 10^4 \text{ m}^3 \text{ s}^{-1}$, estimated from a typical tidal prism of $9 \times 10^8 \text{ m}^3$ in a 12-h tidal cycle (Table 1). The prevailing wind direction is southwesterly, especially in fall and winter, as revealed in the weather model data by the Royal Dutch Meteorological Institute (KNMI, Fig. 2) and by historical observations (Mulder & Louters, 1994). However, occasionally, northeasterly winds are more frequent, for example, April to June in 2009 (Fig. 2b). Wind speeds are generally less than 10 m/s, but the frequency of strong winds ($> 15 \text{ m s}^{-1}$) increases in fall and winter (Fig. 2).

3 Methods

3.1 Model description

The flushing timescale is estimated using the open-source General Estuarine Transport Model (GETM, <https://getm.eu>), a finite-difference hydrostatic primitive-equation model that has been widely applied in estuaries and coastal bays in water exchange studies (e.g., Banas & Hickey, 2005; Duran-Matute et al., 2014). GETM was run on a $300 \text{ m} \times 300 \text{ m}$ Cartesian grid with 10 uniform sigma layers for years 2009–2010. The model domain covers the Oosterschelde and part of the adjacent North Sea (Fig. 1) with a total of 237×111 grid cells. Our model solves flooding and drying of intertidal flats and 3D temperature, salinity, and current velocity forced with air-sea and open boundary conditions. The $k\text{--}\epsilon$ turbulence closure scheme was computed in the 1D General Ocean Turbulence Model (GOTM) coupled with GETM. For the advection scheme of velocity and scalars (temperature, salinity, and tracer), we used the third-order total variation diminishing TVD-P2-PDM and second-order TVD-SUPERBEE for horizontal and vertical dimensions, respectively, which are relatively computationally efficient (Banas & Hickey, 2005).

Bathymetry of the Oosterschelde was interpolated from the regular measurements (resolution 20–200 m) by the Dutch government agency Rijkswaterstaat (accessible from <http://opendap.deltares.nl/thredds/catalog/opendap/hydrografie/surveys/catalog.html>).

Atmospheric forcing including hourly mean winds, air pressure, air temperature, precipitation, and humidity was acquired from the KNMI model results (the downscaled 2.5-km resolution weather model HARMONIE). GETM uses the bulk formulae (Kondo, 1975) to compute the air-sea heat and momentum flux. We used the Flather boundary condition and have prescribed water level and 2D velocity on the open boundary, which were obtained from the output of the Northwest European Shelf tidal prediction and assimilation by Oregon State University (Egbert et al., 2010). Note that this tidal model does not include variations of sea surface height induced by remote wind forcing. The open-boundary temperature and salinity were provided by a 5-km-resolution North Sea model (van der Molen et al., 2016). Freshwater discharge into the northern branch (Fig. 1) was accessed from the Rijkswaterstaat website (<https://waterinfo.rws.nl>). The model displayed acceptable skills in reproducing realistic water level, current velocity, temperature, and salinity, and we refer to an earlier paper for details in detailed model setting, calibration and validation (Jiang et al., 2019a).

3.2 Tracer experiment and calculation of flushing timescales

The conservative Eulerian tracer was placed inside the Oosterschelde to estimate the flushing timescale based on its decay rate from the basin. We define the storm surge barrier as the natural boundary of the Oosterschelde, and the tracer was initialized at unity concentration in every grid cell inside the barrier, and zero outside.

There are many ways to quantify spatially-varying flushing timescales based on tracer decay. Residence time (T_r) of every grid cell could be computed by integration of the remnant function (Takeoka, 1984), $T_r = \int_0^\infty C(t)/C_0 dt$, where $C(t)$ and C_0 are the instantaneous and initial tracer concentration in each grid cell. Since flushing of the tracer inside the domain is never complete, a sufficiently long integration period (duration of the tracer experiment) needs to be arbitrarily defined (e.g., Wang et al., 2004; Yuan et al., 2007), which results in numerical inaccuracy and computational burden, especially for systems/sub-regions with long residence times (Viero & Defina, 2016). An alternative timescale, turnover time (T_t) as defined by Zimmerman (1976), is the time when $C(t) = e^{-1} C_0$. To estimate T_t , a much shorter model run is needed compared to T_r , and it does not need post-process integration for each grid cell, making it computationally efficient. Assuming that the tracer concentration decreases exponentially, i.e., $C(t) = C_0 \cdot \exp(-t/T_t)$, T_t and T_r are equivalent.

The exponential decay assumption holds when the system acts as a continuously stirred tank reactor (Monsen et al., 2002), which should be the case for each grid cell (Ahmed et al., 2017). In tide-dominated bays with negligible river discharge, T_t is close to T_r and suitable for representing the local flushing capability (Rayson et al., 2016; Viero & Defina, 2016). The Oosterschelde is such a tide-dominated and well-mixed system (Section 2), in which the tracer decay is quasi-exponential in all compartments, and T_t and T_r show similar magnitude and spatial patterns (Jiang et al., 2019a). Besides turnover time, T_t has been named flushing (Grifoll et al., 2013; Plus et al., 2009; Sandery & Kämpf, 2007), renewal (Koutitonsky et al., 2004), residence (Abdelrhman, 2005), and influence time (Delhez et al., 2014).

After each model run, time series of the tracer concentration in each grid cell went through a 48-hour low-pass filter to remove the daily tidal signal, and the local T_t is defined as the time when the filtered concentration fell to e^{-1} of the initial concentration. T_t in the entire basin and its four main compartments (Fig. 1) was derived by finding the e^{-1} threshold in the decay curve of the basin- and compartment-wide integrated mass.

3.3. Intra-annual scenarios and sensitivity experiments

In order to examine the intra-annual variability of turnover times in the Oosterschelde, tracer experiments were started at four different times in 2009, 1 January, 1 April, 1 July, and 1 October, following the method of Luff & Pohlmann (1995) and Sandery & Kämpf (2007). Each run lasted six months, slightly longer than the maximal T_t in all grid cells. These intra-annual scenarios were conducted under realistic forcing (R1, R4, R7, and R10 in Table 1). GETM allows decomposing the role of individual forcing agents by switching off others. Such sensitivity tests were performed to examine the relative importance of tides, winds, and gravitational force to the intra-annual variability of turnover times (Table 1). Given that the minimal freshwater flow from the northern branch lacks seasonality for being manipulated by sluice opening and only affects T_t and salinity locally. Based on our prior examination of including and excluding freshwater discharge, the difference in T_t is less than one day in the entire bay and its western, central, and eastern compartments. Thus, it is switched off in all the following sensitivity tests.

In this study, since the tidal phase vary in different parts of the basin (e.g., the lag of high tide is ~30 min from the mouth to the head), the system-wide tidal phase is defined using the basin volume (Duran-Matute et al., 2014). Specifically, the time between two consecutive peaks/troughs of the Oosterschelde volume is considered as one tide cycle. When releasing tracer at five phases during one tidal cycle on 1 January 2009 (STP1–STP5 in Table 1), turnover time T_t exhibited a dependence on the initial basin volume, but not on the direction of tidal currents. Experiments starting on high tides (a larger initial volume of tracer) generated T_t ~10 days longer than that starting on low tides (a smaller initial volume of tracer, Table 1). To exclude this effect from runs examining other forcing mechanisms, the initial basin volume of subsequent experiments was prescribed as 2.82 km³, the same as that at 9:00 on 1 January 2009 (STP2 in Table 1) and close to that at zero water level (Nienhuis & Smaal, 1994). The tide-only STP2 case is the baseline scenario for subsequent runs; i.e., in any sensitivity runs, only one forcing agent was changed compared to STP2 (Table 1).

Since tides predominantly drive flushing of the Oosterschelde, we also studied the intra-annual variability of T_t in experiments driven only by tides in different times of year were examined (STP2, STA4, STA7, and STA10 in Table 1), including two cases starting at spring and neap tides (STAsp and STAne in Table 1).

Wind impacts were assessed by imposing realistic winds in different times of year to the baseline scenario (W1, W4, W7, and W10 in Table 1). Since winds are highly variable with the period of a few days, we ran idealized experiments of fixed prevailing wind directions (southwesterly and northeasterly, Fig. 2) to better understand the wind-modulated flushing processes. Inspired by the method of Scully (2010), a time series was created where the wind magnitude changed sinusoidally from zero to a maximum speed with a period of three days. In accordance with the realistic wind distribution (Fig. 2), we used the maximum speed of 5 m/s, 10 m/s, and 15 m/s to represent weak, intermediate, and strong southwesterly winds (Wsw5, Wsw10, and Wsw15 in Table 1). One scenario (maximum 10 m/s, Wne10 in Table 1) was run for the less prevailing northeasterly winds.

The importance of density-driven (gravitational) flow varies among systems and different regions of a system (Babson et al., 2006; Choi & Lee, 2004; Du et al., 2018) and is largely unknown in the Oosterschelde. The gravitational force induced by seasonal density variations was imposed to the baseline run to demonstrate the baroclinic impacts on the flushing timescale (G1, G4, G7, and G10 in Table 1). In this study, sigma-t, a unitless value equivalent to seawater density minus 1000 kg m⁻³, is used to denote seawater density (Fig. 3). Because of the temporally varying nature of density (Fig. 3), it is not easy to establish the

direct connections between density gradient and T_t . Thus, we set up idealized “fixed-density” baroclinic runs in which density is determined by the initial condition and does not change with time. Events with low and high axial density gradients were selected during the year (Fig. 3) to feed the “fixed-density” runs (GF1–GF4 in Table 1).

4. Intra-annual variability of turnover times

The T_t distribution in the realistic scenarios exhibited a west-east gradient with the longest T_t appearing in the landward end of the eastern compartment (Figs. 4 and 5a). Along the x axis between 35 and 55 km, T_t was significantly shorter in the south of the basin than in the north (Fig. 5a), which indicates that the southern channel acts as the main conduit for tidal flushing. T_t in the deep channels tended to be shorter than that in the surrounding shallow regions; for example, in the eastern compartment, the areas with relatively short T_t marked the location of tidal channels (Fig. 4). These patterns in flushing timescales are also reported in other tidal systems, described as “tongue-shaped” distribution by Shen & Wang (2007).

Table 1 shows T_t of the entire bay and the least flushed eastern compartment, to indicate and intercompare the average and upper limit of the flushing timescale in all presented scenarios. When starting tracer experiments in different seasons, the basin-wide T_t varies no longer than 10 days. However, significant intra-annual differences are manifested in the T_t spatial distribution (Figs. 4 and 5). Directions and spacing of T_t contours were highly variable in the central compartment (Fig. 4). The north of the western compartment is characterized by a large spatial T_t gradient, as revealed by, for example, a 60-day T_t variation within 5 km (Fig. 4a), as well as strong intra-annual variability (Fig. 5b). Besides this area, the intra-annual variability of T_t was pronounced along the southern channel, and in the east of the eastern compartment (Fig. 5b). The longest T_t of the basin ranges from 141 days (Fig. 5d) to 166 days (Fig. 5c). Overall, shorter T_t was found in scenarios starting on 1 July (Fig. 4). The considerable intra-annual variability of T_t implies that the Oosterschelde driven by realistic forcing is not at a steady state. In what follows, the contribution of individual forces is examined.

4.1. The effects of tides on the intra-annual variability of turnover times

As mentioned in Section 3.3, T_t in the entire bay and four compartments differed up to 10 days when starting at various phases of a tidal cycle (STP1–STP5 in Table 1). With the similar initial basin volume, tracer experiments starting at flood (15:00) and ebb (09:00) tides generated similar T_t (STP2 and STP4 in Table 1), which reveals that the initial basin volume, rather than the direction of tidal currents, is the major factor causing the T_t variability when starting at different tidal phases.

The importance of tidal amplitude is shown by the fact that the residence time was almost doubled owing to the declined tidal amplitude after construction of the storm surge barrier (Nienhuis & Smaal, 1994; Wetsteyn & Kromkamp, 1994). However, over a time span comparable to T_t (~150 days), with the same starting tidal phase and basin volume (Section 3.3), the average realistic tidal amplitude is relatively invariable; i.e., based on our calculation, the average tidal amplitude during Days 1–150 is nearly the same as that during Days 91–240, 121–270, 151–300, etc (standard deviation 2.8 mm). Thus, in the tide-only scenarios with realistic tides starting at different times of year (STP2, STA4, STA7, and STA10 in Table 1), T_t did not vary substantially (Table 1). The basin-wide and eastern-compartment turnover times estimated in the scenario starting at the spring tide are around two days shorter than those in the scenario starting at the neap tide (STAsp, and STAne in Table 1), which, nevertheless, is insufficient to explain the intra-annual variability in the realistic simulations. These findings indicate that tides are not a significant source of intra-annual variability.

4.2. The effects of winds on the intra-annual variability of turnover times

The prevailing southwesterly winds varied in frequency and magnitude within the year 2009 (Fig. 2). Compared to the tide-only scenario, imposing winds did not exert a large change in the average T_t of the entire bay but reduced the T_t in the eastern compartment by 2–14 days (Table 1). Spatially, similar wind-induced reduction in T_t was manifested in many shallow regions including the north of the western compartment, southeast of the eastern compartment, and most of the northern branch (Figs. 6a–6d). Despite these similarities, the spatial patterns of T_t differed due to the wind intra-annual variability (Figs. 6a–6d). Turnover times in these shallow regions were shorter in scenarios with winds starting in July and October (Figs. 6a–6d). T_t in the central compartment in the scenario with winds starting on 1 October increased compared to the baseline scenario, which was not the case in the other three scenarios (Figs. 6a–6d). Given the high temporal variability of realistic winds, idealized winds were used to explore the wind-modulated spatial patterns in T_t and understand the above findings.

Under southwesterly winds of maximum 5 m/s, the wind-driven changes in T_t were limited to the eastern and northern compartments (Fig. 6e). When southwesterly winds were increased to maximum 10 m/s and 15 m/s, the spatial pattern of T_t resembled that driven by the realistic winds starting on 1 October (Figs. 6d, 6f, and 6g). Thus, southwesterly winds tended to increase T_t in the central compartment and decrease it in many shallow (mostly < 5 m) areas of other compartments, and these effects strengthened with wind magnitude (Figs. 6e–6g and 7). In contrast, under the less prevailing northeasterly winds of maximum 10 m/s, T_t was reduced broadly in the eastern, northern, and central compartments and basin-wide (Fig. 6h and Table 1). These phenomena could be ascribed to the wind-driven circulation. In the residual current field driven by southwesterly winds, waters from the surrounding compartments were transported toward the central compartment, where a convergent anticlockwise gyre was generated (Fig. 8a) and tracer accumulated (Fig. 8c). Under the northeasterly winds, the residual currents were reversed and featured an outward flow (Fig. 8b), diminishing the tracer concentration in most of the basin (Fig. 8d).

It should be noted that, although idealized wind scenarios cannot be used directly for T_t estimations, they do offer insight into the spatial patterns in the intra-annual scenarios. Southwesterly winds are frequent in all seasons (Fig. 2), so all scenarios driven by realistic winds demonstrate the typical patterns of southwesterly winds, that is, significant T_t declines in most shallow waters (e.g., up to 5 m deep, Fig. 7). Due to the prevalence of strong (> 10 m/s) southwesterly winds, increased T_t was found in the central compartment in the scenario with winds starting on 1 October, a feature induced only by strong southwesterly winds (Figs. 2, 6d, 6f, and 6g). In summary, it is found that winds of sufficient strength (> 5 m/s) greatly altered the spatial heterogeneity of flushing characteristics. Although the average T_t of the entire basin was not greatly changed by winds (Table 1), wind-induced changes in T_t were as high as 40 days in certain regions (Fig. 6).

4.3. The effects of gravitational flow on the intra-annual variability of turnover times

The spatial density gradient in the Oosterschelde is shaped by temperature and salinity. The annual cycle of temperature causes density to be lower in summer and higher in winter (Fig. 3). However, it is salinity that has the main impact on spatial density variations. During freshet events of the Westerschelde and Rhine rivers, salinity at the southern and northern open boundaries can be lower than that in the Oosterschelde (e.g., Fig. 9a), reducing axial density gradient between the head and mouth of the basin (Fig. 3). When these two river plumes are so small that salinity near the mouth of the Oosterschelde is not greatly affected (e.g., Fig. 9b), the bay is characterized by a larger density gradient in the east-west direction

(Fig. 3). Therefore, the salinity out of the mouth of the Oosterschelde is regulated by mixing of the North Sea water and outflow of the two rivers, and discharge and plume morphology of the two rivers seem to affect the horizontal density gradient in the Oosterschelde (Figs. 3 and 9).

The gravitational force induced by such density variations created substantial differences in the T_t magnitude and distribution (Table 1 and Fig. 10). Because the density-driven exchange flow favors flushing of the basin, T_t was shortened in all four scenarios including it (Table 1 and Figs. 10a–10d). However, the extent of T_t reduction varied dramatically among scenarios, and flushing was accelerated the most and least in Figs. 10c and 10a, respectively (see also Table 1). Given the density gradient dynamically changes within the duration of a scenario (6 months, Fig. 3), idealized scenarios with “fixed” density were examined to establish the relationship between density and T_t . When density was the only variable among scenarios and stationary with time, it is found that a relatively sharp axial density gradient promotes flushing by driving a strong bidirectional gravitational flow (Figs. 10e, 10f, and 11). The relationship between the axial density gradient, exchange flow, and T_t was found to hold in the scenarios with realistic gravitational flow (Table 1 and Figs. 10a–10d and 12). Note that density and current results in Fig. 12 are 90-day averages to avoid overlaps of scenarios. Therefore, we conclude that intra-annual density variations exert a significant influence on flushing of the Oosterschelde.

5. Discussion

Of various methods used to evaluate the local flushing timescale of a tidal bay, our approach estimating turnover times with Eulerian tracer exhibits certain advantages. When examined at different tidal phases, results of our tracer experiments were insensitive to the initial direction of tidal currents (Table 1). This is different from the Lagrangian particle-tracking method, in which the phase of the initial tidal current is highly influential on movement of every single particle (Monsen et al., 2002; Oliveira & Baptista, 1997). Our usage of turnover time, similar to the magnitude of residence time in tide-dominated systems (Jiang et al., 2019a; Viero & Defina, 2016), is relatively computationally efficient, making it possible to run realistic and idealized sensitivity experiments decomposing the intra-annual variability. Approaches in this study are recommended to decipher the spatiotemporal variability of flushing processes in other tidal bays.

Of all the forces modulating flushing of the Oosterschelde, tides predominately drive the water renewal in the Oosterschelde, and flushing timescales would be indefinitely long without them. Effects of winds and gravitational flow on water renewal were manifested by modifying the tidal currents (e.g., Figs. 8 and 11). Tidal variations on daily (diurnal and semi-diurnal) and biweekly (spring-neap) scales are found important to systems with flushing timescales of days to weeks, such as the estuaries of Richibucto (Guyondet et al., 2005), Danshuei River (Wang et al., 2004), Yaquina Bay (Lemagie & Lerczak, 2015), Mersey (Yuan et al., 2007), Suances (Bárcena et al., 2012), Willapa Bay (Banas & Hickey, 2005). However, the trimonthly and six-monthly average tidal amplitude were almost invariable (Section 4.1). The fact that starting the sensitivity experiments with tides of different seasons but the same phase (i.e., basin volume) resulted in insignificant changes in turnover times reveal that tidal and subtidal processes are not the dominant contributor to the T_t in the Oosterschelde over the multi-month scale.

Compared to tides, the impacts of winds and gravitational flow are much more important in creating the intra-annual variability of turnover times in our system. Winds were responsible for altering the residual currents and redistributing the tracer concentration in the basin (Fig. 8), which induced different spatial patterns of turnover times. This mechanism

resembles findings in the Venice Lagoon (Cucco & Umgiesser, 2006). Shallow waters, where relatively weak winds can affect the local hydrodynamics, are particularly susceptible to wind forcing, as found in our results and previously (Geyer, 1997; Ranasinghe & Pattiaratchi, 1998). In estuaries, wind-induced mixing can also modulate the density field and thus impact flushing (e.g., Du et al., 20018); however, this effect is not so important here according to our examination since tides are usually strong enough to mix the water column (Mulder & Louters, 1994).

As a tidal bay, the horizontal density gradient in our system ($\Delta T < 3\text{ }^{\circ}\text{C}$ and $\Delta S < 3$ based on our 2-year simulation) is not as strong as in many river-dominated stratified estuaries such as Chesapeake and Mobile Bays, where gravitational flow is vital for flushing of the basin (Du et al., 2018; Shen & Wang, 2007). In spite of the small density gradient in our study, gravitational exchange flow could reduce the flushing timescale substantially (> 20 days) and vastly (Fig. 10). Thus, studies neglecting the gravitational force or vertical dynamics may lead to overestimates of flushing timescales, especially in systems with important buoyancy sources (e.g., Bárcena et al., 2012; Guyonnet et al. 2005; Webb & Marr, 2016).

In our model, tides are prescribed by imposing water elevation at the open boundary, which is optimized by data assimilation and validation of the shallow-water tides using tidal gauge data (Egbert et al., 2010). It should be noted that this open boundary water elevation, despite producing a good fit with the observed Oosterschelde water level in our model (Jiang et al., 2019a), does not include the subtidal effects on coastal sea surface height, which is mostly driven by remote winds (Frederikse and Gerkema, 2018). If the impacts of remote wind forcing are considered, the variability of turnover times should be enhanced in winter and fall since these two seasons are found to exhibit the largest wind-driven sea level variability along the North Sea coast (Frederikse and Gerkema, 2018).

6. Summary and conclusions

In this study, tracer experiments using realistic hydrodynamic model runs were used to investigate the intra-annual and spatial variability of turnover times in a tidal bay, the Oosterschelde. Our study displays the strong intra-annual spatial variability of turnover times in the Oosterschelde and its interactions with the basin geometry and bathymetry and sensitivity to various forcing. In such a dynamic system, spatially mapping of T_t offers more information than a single bulk/integrative timescale for the whole basin or sub-region. For example, the basin-wide bulk turnover time was similar (85.9–89.2 days) among scenarios with tides and realistic winds (W1, W4, W7, and W10 in Table 1), whereas the spatial distribution varied strongly (Fig. 6a–6d).

The mechanisms behind the intra-annual variability were studied based on numerical experiments isolating certain processes while eliminating others. Our results indicate that turnover times in the Oosterschelde are sensitive to the initial basin volume but not to the starting tidal current directions. When starting the tracer experiments with the same initial volume (mass) of tracer, the starting tidal amplitude (spring vs neap) is also not a major source of the intra-annual variability of T_t . Compared to tides, winds and gravitational flow were found to be predominant in shaping the intra-annual and spatial patterns in T_t . For example, strong southwesterly winds were responsible for the relatively long T_t in the central compartment in the intra-annual scenario starting on 1 October (Figs. 2 and 4); stronger density gradients along the main stem accounted largely for the overall faster flushing in the intra-annual scenarios starting on 1 July (Figs. 4 and 11). The role of gravitational force in promoting the basin flushing is further corroborated by “fixed-density” scenarios in which the density gradient is constant during experiments.

Our results highlight the role of surrounding rivers, by acting as external buoyancy sources near the bay mouth, in mediating salinity and flushing of the Oosterschelde. In order to protect the shoreline from flooding and erosion, the southwestern delta in the Netherlands is highly manipulated by construction of dams and sluices, and opening of previously closed surrounding systems (e.g. Grevelingen and Haringvliet) is underway (van Haren, 2019; Ysebaert et al., 2016). Findings in this study imply that these engineering works should take the interaction of basins into account instead of focusing on a single system. Salinity changes in the coastal ocean can potentially change the flushing characteristics and ecosystem functions of each basin, since, for example, the Oosterschelde and Grevelingen, both used extensively for shellfish culture, can benefit from enhanced tidal import of seston from the North Sea (Smaal et al., 2013). Hence, continued oceanographic monitoring and numerical modeling covering the whole delta region are essential to understand the inter-basin dynamics and assisting environmental managers in decision-making.

A systematic examination of each forcing improves our understanding of flushing processes in a tidal bay. Compared to previous works, our results advance the efforts to interpret the mechanisms of wind- and density-driven circulation in this region. Here we examined the case of the Oosterschelde, but the approaches described here can be applied more generally to other tidal bays.

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640 **Table 1.** A list of model runs with tracer experiments and the resulting turnover times (T_t) in the year 2009. Five tidal phases in STP1–STP5
641 represent high tide, ebb, low tide, flood, high tide, respectively. In STAsp and STAne scenarios, it is spring and neap tide on the starting date 13
642 Jan. and 20 Jan. The density is “fixed” (static) in the last four runs GF1–GF4; namely, the gravitational flow is determined by the initial density
643 gradient. Model data are archived on an online dataset (Jiang et al., 2019b).

Run name	Startup time	Initial basin volume (km ³)	Startup Temperature / Salinity	Tides	Winds	Density calculation	T_t (days) basin-wide	T_t (days) eastern compartment
Realistic intra-annual runs								
R1	1 Jan. 0:00	2.68	Realistic	Realistic	Realistic	Baroclinic	81.7	137.1
R4	1 Apr. 0:00	2.48	Realistic	Realistic	Realistic	Baroclinic	74.6	134.5
R7	1 Jul. 0:00	2.88	Realistic	Realistic	Realistic	Baroclinic	79.8	129.5
R10	1 Oct. 0:00	3.15	Realistic	Realistic	Realistic	Baroclinic	81.5	136.8
Sensitivity runs of the startup tidal phase								
STP1	1 Jan. 06:00	3.37	No	Realistic	No	Barotropic	92.9	153.5
STP2	1 Jan. 09:00	2.82	No	Realistic	No	Barotropic	87.2	147.7
STP3	1 Jan. 12:00	2.39	No	Realistic	No	Barotropic	83.7	143.5
STP4	1 Jan. 15:00	2.59	No	Realistic	No	Barotropic	85.3	145.4
STP5	1 Jan. 18:00	3.29	No	Realistic	No	Barotropic	92.4	152.9
Sensitivity runs of the startup tidal amplitude								
STA4	1 Jan. 09:00	2.82	No	Realistic starting on 1 Apr.	No	Barotropic	88.6	148.3
STA7	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jul.	No	Barotropic	89.3	149.9
STA10	1 Jan. 09:00	2.82	No	Realistic starting on 1 Oct.	No	Barotropic	87.9	149.3
STAsp	1 Jan. 09:00	2.82	No	Realistic starting on 13 Jan.	No	Barotropic	87.0	148.3
STAne	1 Jan. 09:00	2.82	No	Realistic starting on 20 Jan.	No	Barotropic	89.1	150.8
Sensitivity runs of winds								
W1	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Realistic	Barotropic	86.9	142.1
W4	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Realistic starting on 1 Apr.	Barotropic	85.9	141.4
W7	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Realistic starting on 1 Jul.	Barotropic	86.2	133.5
W10	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Realistic starting on 1 Oct.	Barotropic	88.2	145.1
Wsw5	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Southwesterly 5 m/s	Barotropic	89.7	148.5
Wsw10	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Southwesterly 10 m/s	Barotropic	95.8	148.3
Wsw15	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Southwesterly 15 m/s	Barotropic	90.2	143.9
Wne10	1 Jan. 09:00	2.82	No	Realistic starting on 1 Jan.	Northeasterly 10 m/s	Barotropic	71.1	116.5

Sensitivity runs of gravitational flow								
G1	1 Jan. 09:00	2.82	Realistic	Realistic starting on 1 Jan.	No	Baroclinic	84.0	143.6
G4	1 Jan. 09:00	2.82	Realistic starting on 1 Apr.	Realistic starting on 1 Jan.	No	Baroclinic	73.7	117.3
G7	1 Jan. 09:00	2.82	Realistic starting on 1 Jul.	Realistic starting on 1 Jan.	No	Baroclinic	66.4	113.7
G10	1 Jan. 09:00	2.82	Realistic starting on 1 Oct.	Realistic starting on 1 Jan.	No	Baroclinic	75.2	129.5
GF1 ^h	1 Jan. 09:00	2.82	25 Feb. “fixed”	Realistic starting on 1 Jan.	No	Baroclinic	85.0	147.0
GF2 ^h	1 Jan. 09:00	2.82	26 Aug. “fixed”	Realistic starting on 1 Jan.	No	Baroclinic	69.5	125.3
GF3 ^h	1 Jan. 09:00	2.82	21 Sep. “fixed”	Realistic starting on 1 Jan.	No	Baroclinic	83.3	146.0
GF4 ^h	1 Jan. 09:00	2.82	21 Dec. “fixed”	Realistic starting on 1 Jan.	No	Baroclinic	67.9	119.2

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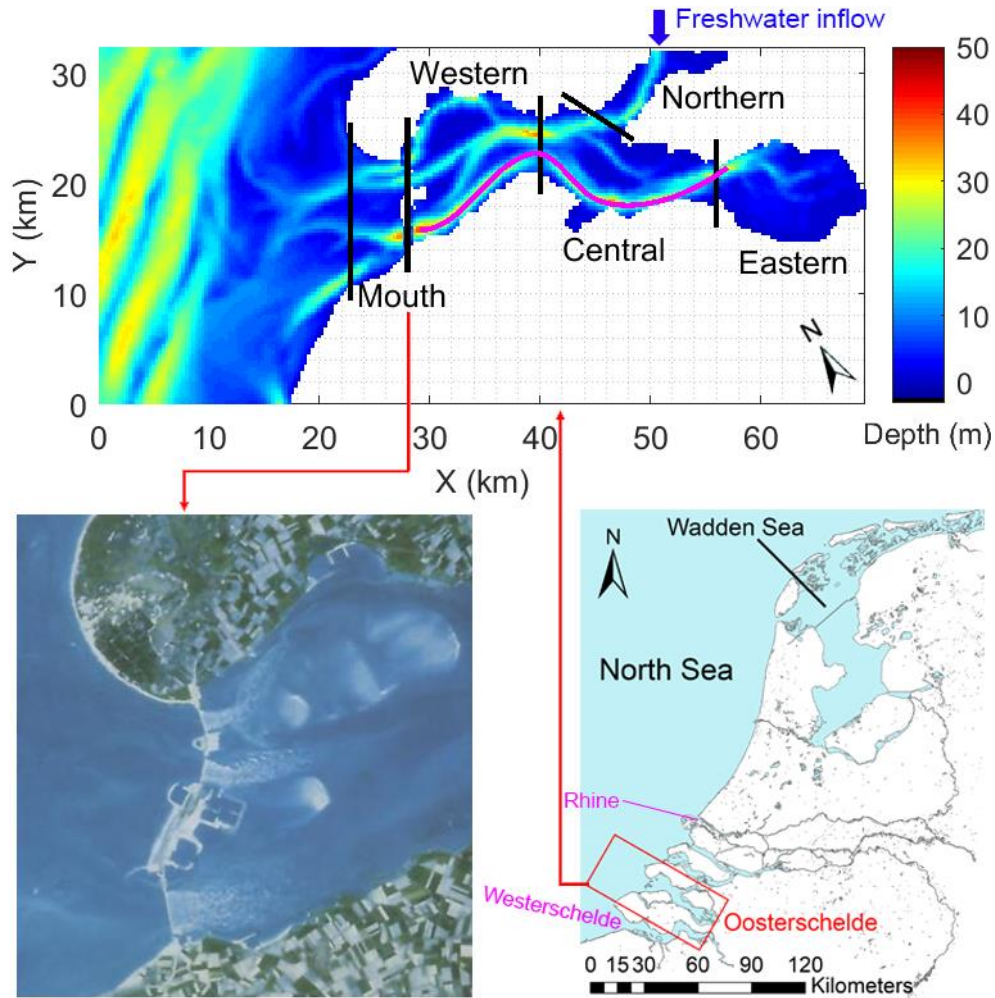


Fig. 1. The model domain and bathymetry covering the Oosterschelde and part of the southern North Sea. The black lines in the left panel divide the Oosterschelde into four compartments (western, central, eastern, and northern) and a “mouth” region. The pink line in the left panel denotes a transact along the main (southern) channel of the basin. An aerial image of the bay mouth (Photograph by National Aerospace Laboratory, Amsterdam, published in Nienhuis and Smaal, 1994) shows three opening of the storm surge barrier.

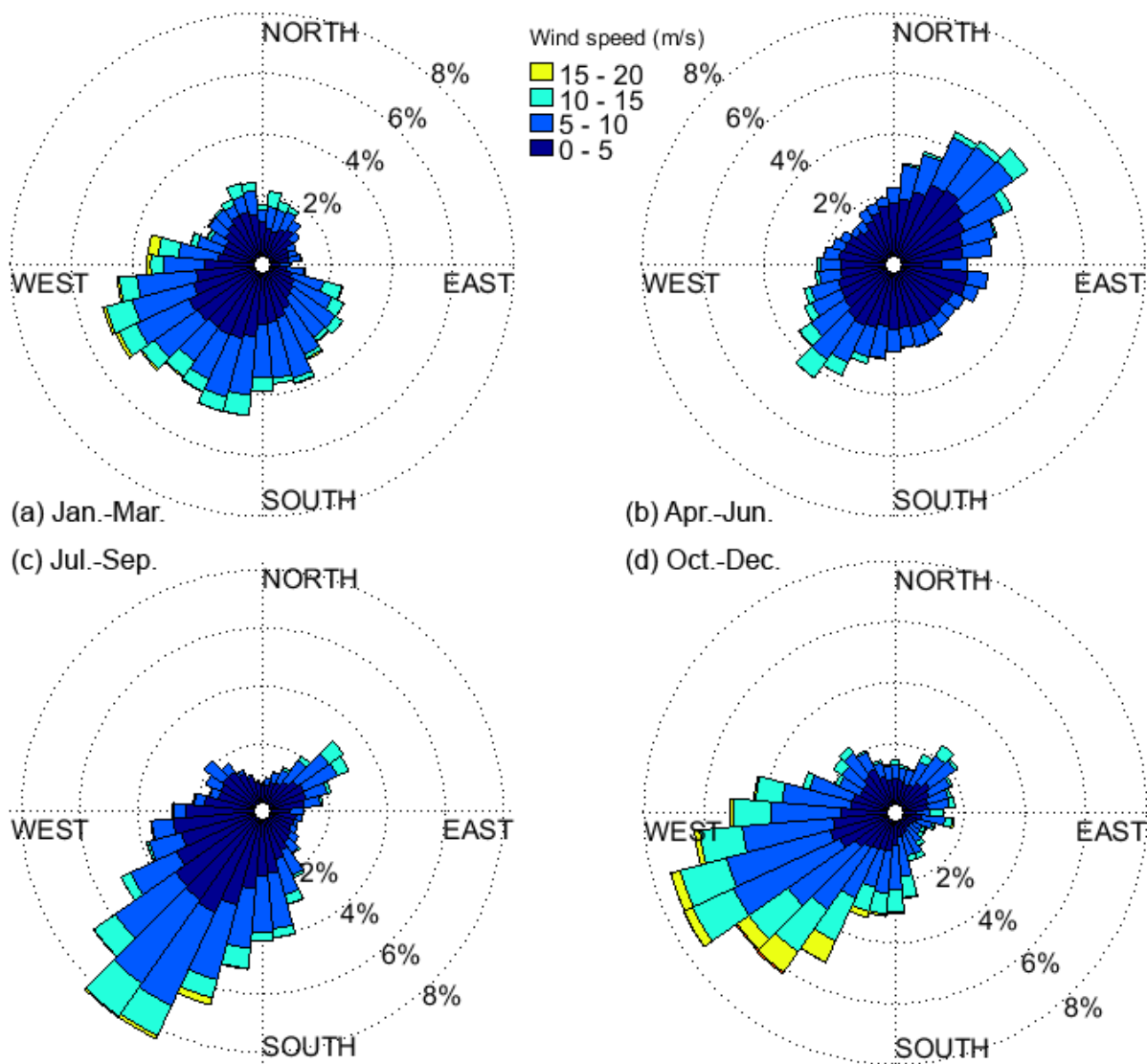


Fig. 2. Trimonthly wind roses over the Oosterschelde in 2009. Wind data are provided by a downscaled weather model HARMONIE with a horizontal grid of 2.5 km produced by the Royal Dutch Meteorological Institute (KNMI).

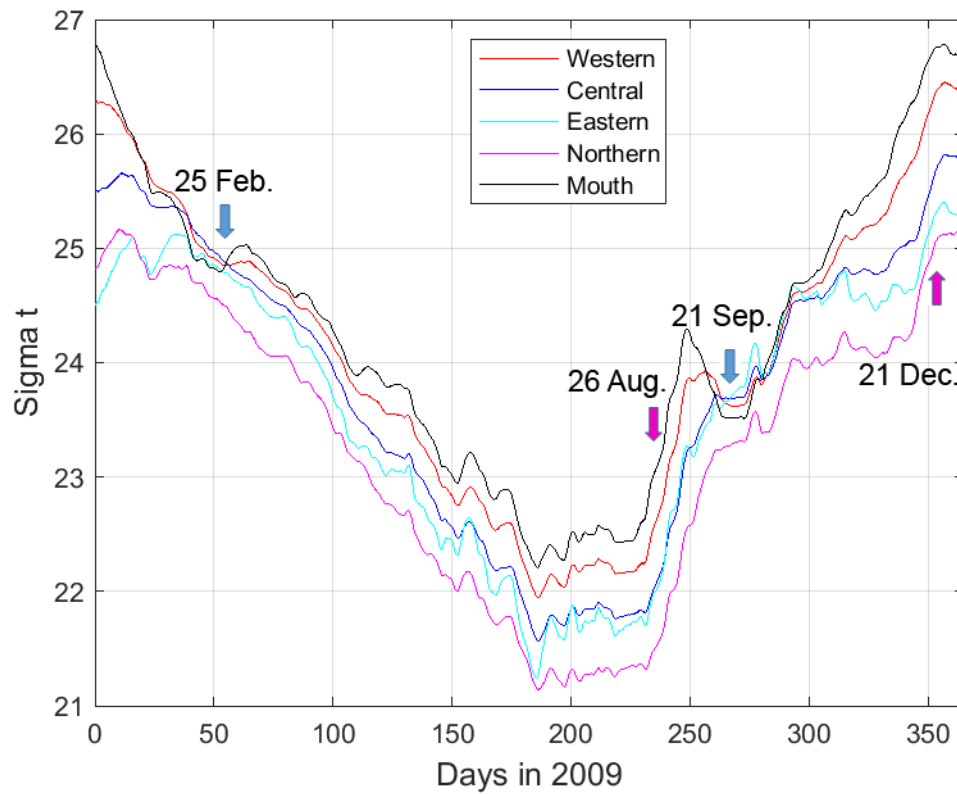


Fig. 3. Simulated spatially average sigma-t in four compartments and the area “mouth” in the Oosterschelde in 2009. Sigma-t is a dimensionless quantity subtracting 1000 kg m^{-3} from the density of seawater, e.g, a sigma-t value of 25 denotes the seawater density 1025 kg m^{-3} . Four dates marked were used for initial conditions of idealized “fixed-density” scenarios (Table 1). See Fig. 1 for the division of these sub-regions.

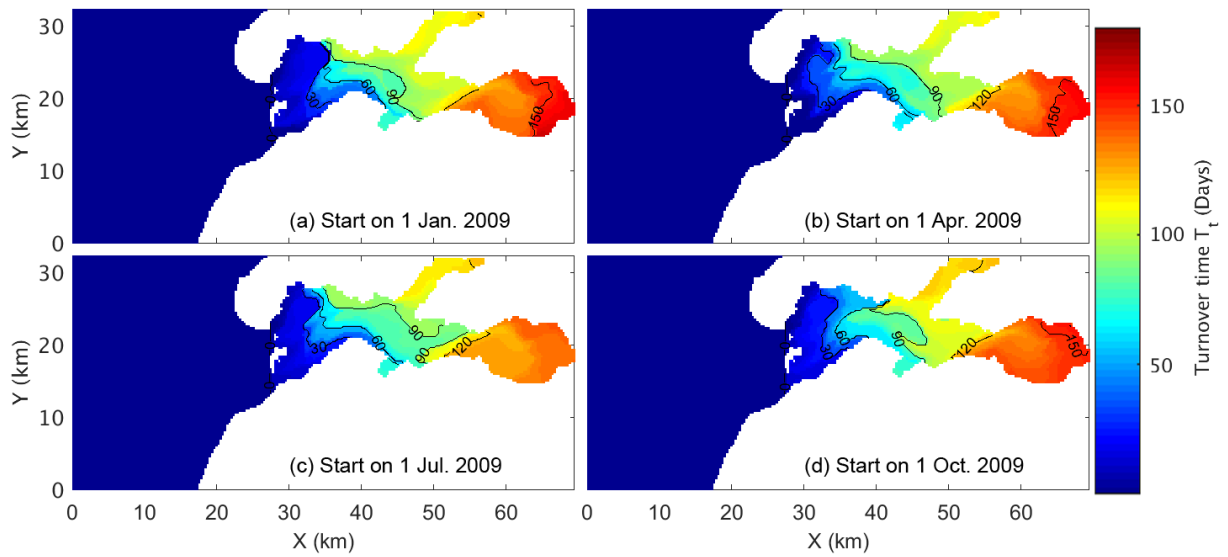


Fig. 4. Turnover time of four realistic scenarios initiating on (a) 1 January, (b) 1 April, (c) 1 July, and (d) 1 October in 2009. See R1, R4, R7, and R10 in Table 1 for details of scenarios.

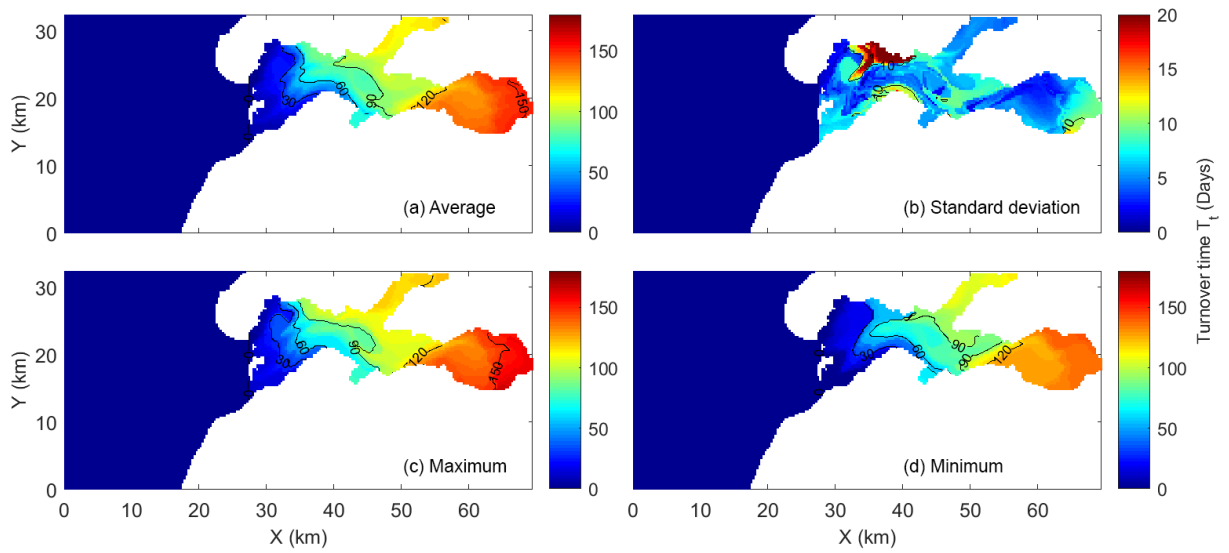


Fig. 5. (a) Average, (b) standard deviation, (c) maximum, and (d) minimum of turnover time in the four scenarios shown in Fig. 4.

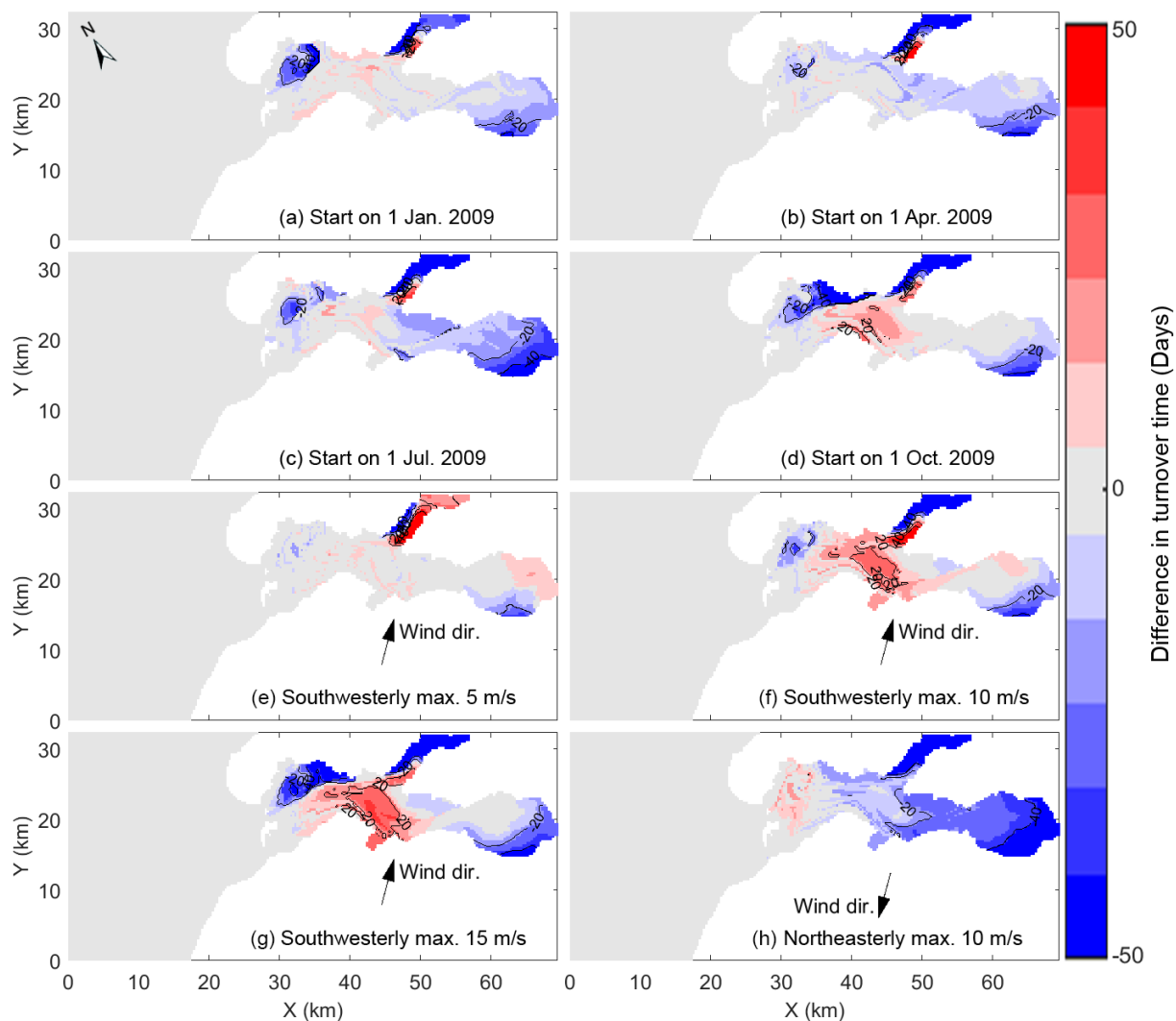


Fig. 6. Differences in turnover time distribution between “tides + winds” and “tide-only” scenarios. The wind forcing is (a-d) realistic and (e-h) idealized with fixed direction: (a) W1 – STP2, (b) W4 – STP2, (c) W7 – STP2, (d) W10 – STP2, (e) Wsw5 – STP2, (f) Wsw10 – STP2, (g) Wsw15 – STP2, and (h) Wne10 – STP2. See Table 1 for details of scenarios.

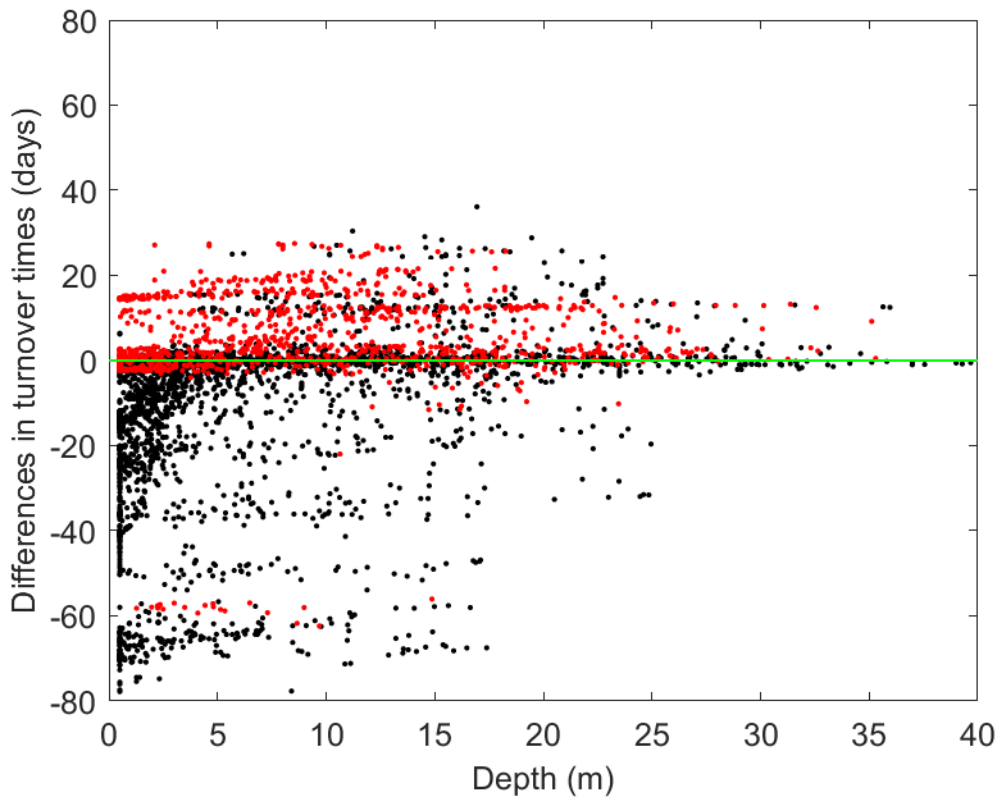


Fig. 7. Differences in turnover time between “tides + winds” and “tide-only” scenarios (Wsw10 – STP2) versus water depth. Red and black points denote each 300 m × 300 m grid cell in the central and the other three compartments, respectively. The green line represents no difference in turnover time.

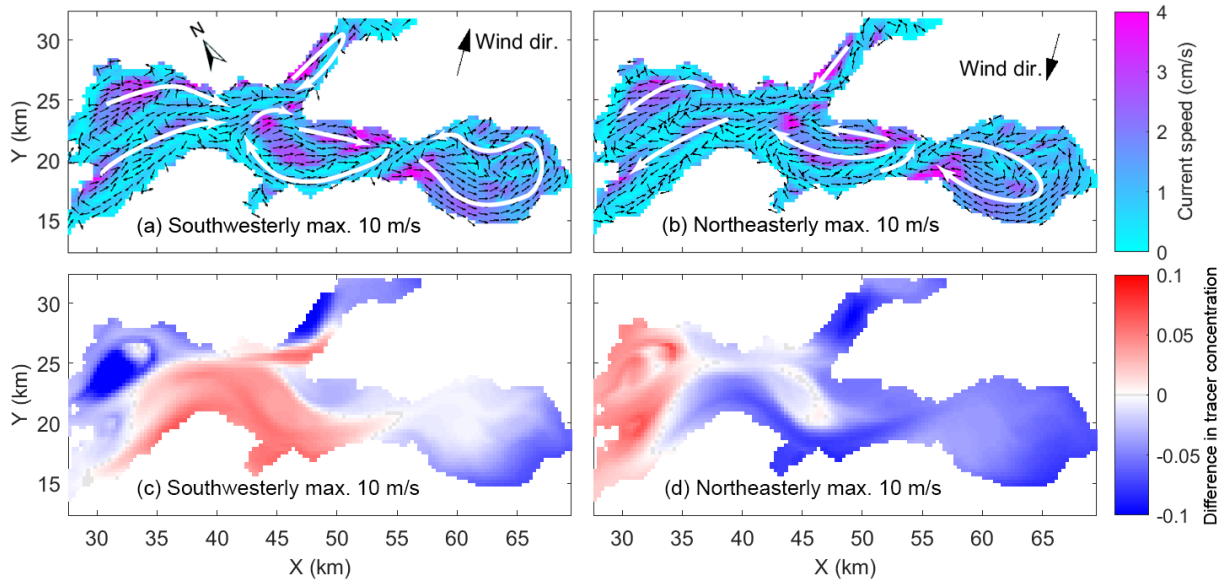


Fig. 8. (a and b) Depth-average residual currents in the first month of the experiment in scenarios Wsw10 and Wne10 with tide and idealized winds. (c and d) Difference of tracer concentration in these two scenarios from the baseline scenario during the same period. See Table 1 for details of the two scenarios. White arrows in (a) and (b) indicate the schematic circulation patterns.

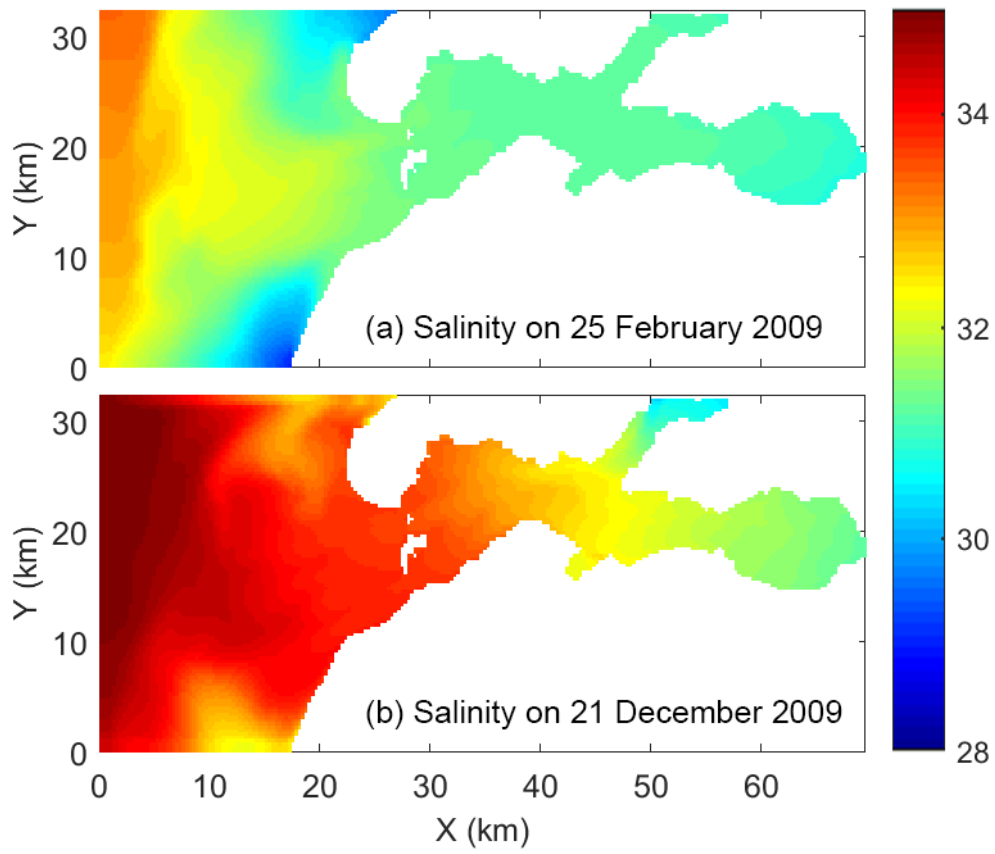


Fig. 9. Depth-average salinity snapshots on (a) 25 February and (b) 21 December 2009. These two dates are marked in Fig. 3.

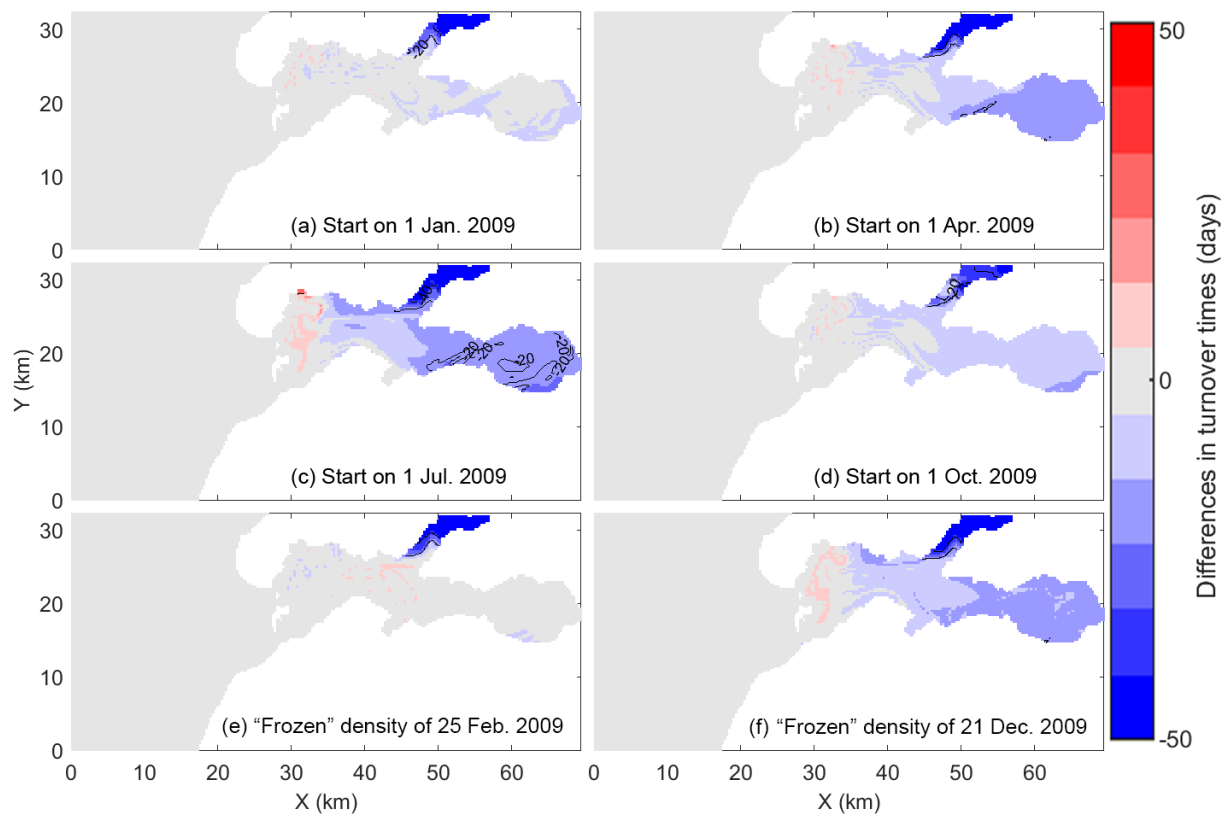


Fig. 10. Differences in turnover time distribution between “tides + gravitational flow” and “tide-only” scenarios. Gravitational flow is (a-d) realistic and induced by (e and f) “fixed” density: (a) G1 – STP2, (b) G4 – STP2, (c) G7 – STP2, (d) G10 – STP2, (e) GF1 – STP2, and (f) GF4 – STP2. See Table 1 for details of scenarios.

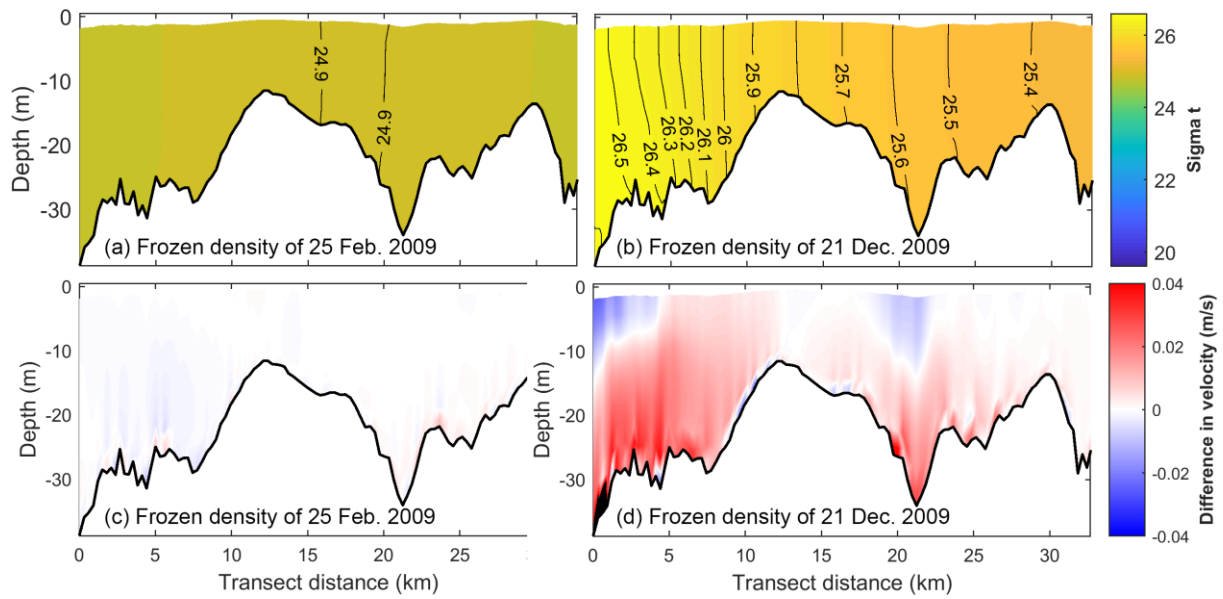


Fig. 11. Density on the transect along the southern channel (Fig. 1) of the Oosterschelde in scenarios (a) GF1 and (b) GF4 with tides and “fixed” gravitational forces. (c and d) Differences in along-channel velocity between “tides + gravitational flow” and “tide-only” scenarios: (c) GF1 – STP2 and (d) GF4 – STP2. See Table 1 for details of scenarios.

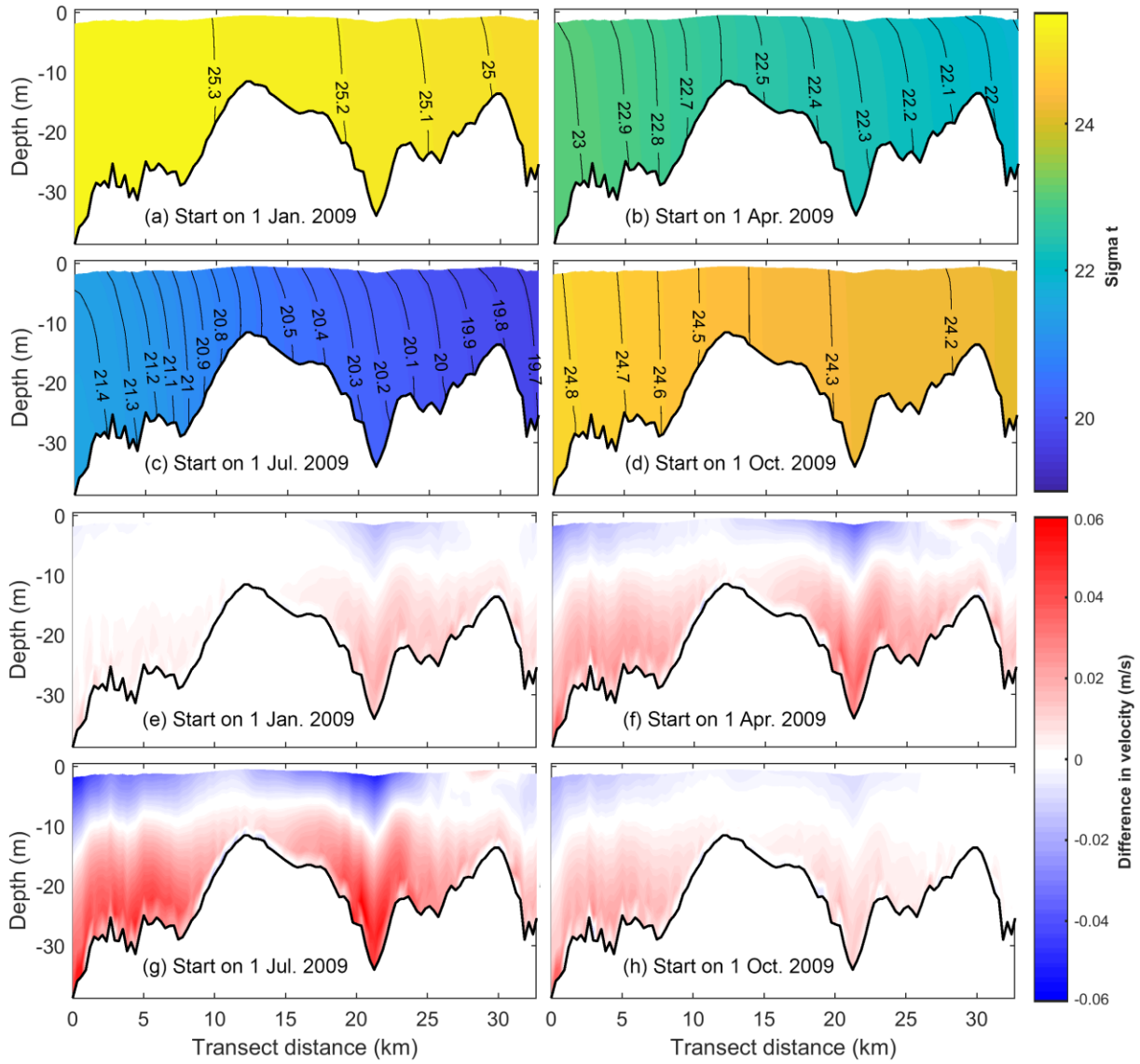


Fig. 12. (a-d) Density on the transect along the southern channel (Fig. 1) of the Oosterschelde in the first 90 days of scenarios G1, G4, G7, and G10 with tides and realistic gravitational forces. (e-h) Differences in along-channel velocity in the first 90 days between “tides + gravitational flow” and “tide-only” scenarios: (e) G1 – STP2, (f) G4 – STP2, (g) G7 – STP2, and (h) G10 – STP2. See Table 1 for details of scenarios.